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INTERPRETATION OF ELECTROMAGNETIC SOUNDINGS  
USING A LAYERED EARTH MODEL

By

Jeffrey J. Daniels

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A Thesis submitted to the Faculty and the Board of Trustees of the Colorado School of Mines in partial fulfillment of the requirements for the degree of Doctor of Philosophy in Geophysical Engineering.

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## ABSTRACT

Theoretical expressions for the electromagnetic field components are computed fast and accurately by using linear filter theory. In this paper a computer program is developed to numerically evaluate the electromagnetic field components ( $E_x$ ,  $E_y$ ,  $E_c$  and  $H_z$ ) for a current dipole source over a layered earth.

Scaled  $H_z$  and  $E_c$  ( $E_c = \sin(z\theta)E_x - \cos(z\theta)E_y$ ) field component curves simplify catalog curve matching by eliminating the curve dependence on  $\theta$ ,  $I$  and  $ds$ . A sample catalog for the  $H_z$  component is given in Appendix A.

A least-squares curve fitting procedure is developed to aid in the interpretation of electromagnetic sounding data. This procedure is tested on two and three layer theoretical sounding curves. Results of these tests can be summarized as follows:

(1) For the two layer case convergence of the layered earth parameters (resistivity and thickness) is accurate to within ten percent.

(2) Three layer curve fitting, with four or more parameters, requires a good first guess to insure convergence. The first guess curve should be within twenty percent of the actual field component values.

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## INTRODUCTION

The objective of controlled-source electromagnetic sounding is to obtain an accurate representation of the subsurface geology. This goal can be achieved by mathematical modeling (Keller and Frischknecht, 1966, Keller, 1968, Frischknecht, 1967, Vanyan, 1965, Pritchard, 1971, Sinha and Bhattacharyya, 1967, Wait, 1966, and others).

Recent advances in numerical techniques have made possible the solution of practical mathematical models. Since these models can only approximate the real geologic situation the interpreter must make use of all auxiliary information at his disposal. This information may come from knowledge of the geology of a given area, well data, or other geophysical data.

One of the more useful models is that for a horizontally layered earth where each layer is represented by an electrical resistivity and a layer thickness. This approach is basically the same as the "apparent resistivity" approach described for electromagnetic sounding data by Vanyan (1967) and for direct-current electrical sounding data by Crous (1971). The layered earth model is applicable to many geological problems (e.g. layering in a sedimentary basement, determination of depths and thicknesses of oil shale and tar sand formations, determination of highly resistant wave guides, etc.). Digital linear filter theory (Ghosh, 1971) provides a rapid technique for computing theoretical electromagnetic sounding curves based on the layered earth model.

Model-based direct interpretation can be achieved by one of the following techniques.

1. Catalog Interpretation - the observed sounding curve is compared to a catalog of theoretical curves (geological) parameters are given by the theoretical curves which matches the observed data),

2. "Cut and Try" Interpretation - model parameters are repeatedly adjusted and new models are generated until the theoretical curve which matches the observed data is found, and

3. Automatic Interpretation - first guess parameters are iteratively adjusted according to some statistical criteria until the theoretical curve matches the observed data. This paper uses a modification of the basic least-squares criterion to develop an automated interpretation scheme.

The purpose of this paper is to develop a computer oriented interpretation technique for electromagnetic sounding data. A layered earth model is assumed to fit the geological conditions represented by the sounding data.

## ELECTROMAGNETIC FIELD COMPONENTS

Figure 1 illustrates the theoretical model considered in this paper. The source is a horizontal current dipole. The model is a horizontally stratified isotropic medium, where  $\rho_i$  and  $h_i$  represent the electrical resistivity and thickness of the  $i$ -th layer. Although the following discussion considers an  $N$ -layered earth the interpretation scheme based on this model assumes a maximum of three layers. The computer programs developed in this paper can easily be

Following Wait's development (1966) the fields vary according to . The free-space wave number is assumed to be equal to zero. The vertical magnetic field component and the horizontal electric field components are expressed as

$$H_z = \frac{I ds}{4\pi} \sin \theta \int_0^{\infty} \lambda (1 + R_1(\lambda)) J_1(\lambda r) d\lambda \quad (2.1)$$

$$E_x = \frac{I ds}{4\pi} \gamma_1^2 \rho_i \int_0^{\infty} (1 + R_1(\lambda)) J_0(\lambda r) d\lambda$$

$$+ \frac{\rho_i I ds}{4\pi} \left\{ -\cos^2 \theta \int_0^{\infty} \lambda^2 \left[ -\frac{\gamma_1^2}{\lambda^2} + \frac{2Z_1(\lambda)}{\lambda \rho_i} - \frac{(\lambda^2 + u_1^2 R_1(\lambda))}{\lambda^2} \right. \right.$$

$$\left. + (1 + R_1(\lambda)) \right] J_0(\lambda r) d\lambda + \frac{\cos(2\theta)}{r} \int_0^{\infty} \lambda \left[ -\frac{\gamma_1^2}{\lambda^2} + \frac{2Z_1(\lambda)}{\lambda \rho_i} \right.$$

$$\left. - \frac{(\lambda^2 + u_1^2 R_1(\lambda))}{\lambda^2} + (1 + R_1(\lambda)) \right] J_1(\lambda r) d\lambda \quad (2.2)$$

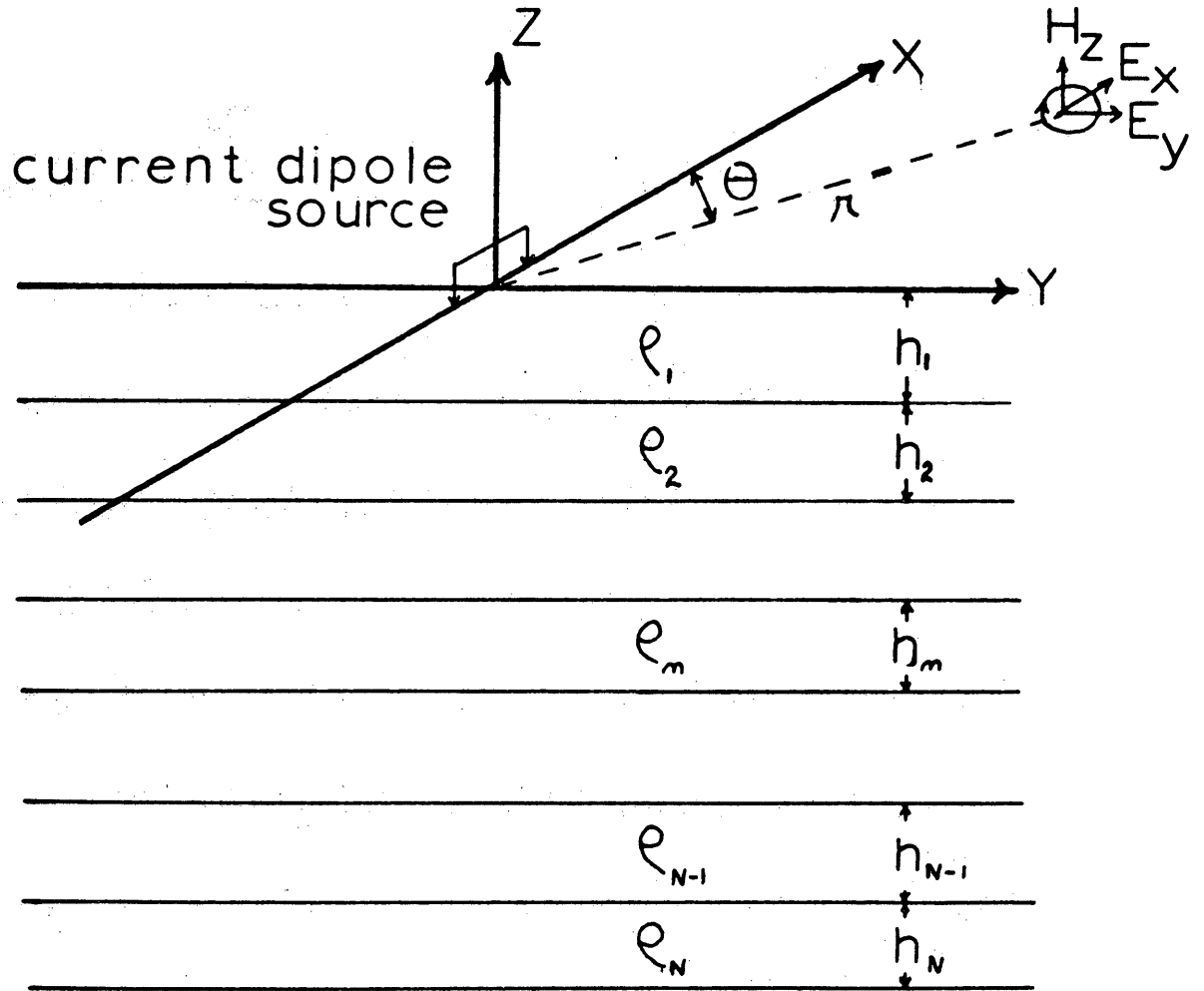


Figure 1. Theoretical model.

$$\begin{aligned}
 \text{and } E_y = \frac{\rho_1 I ds}{4\pi} \left\{ -\frac{\sin(2\theta)}{2} \int_0^\infty \lambda^2 \left[ -\frac{Y_1^2}{\lambda^2} + \frac{2Z_1(\lambda)}{\lambda \rho_1} \right. \right. \\
 \left. \left. - \frac{(\lambda^2 + u_1^2 R_\perp(\lambda))}{\lambda^2} + (1 + R_\perp(\lambda)) \right] J_0(\lambda r) d\lambda + \frac{\sin(2\theta)}{\pi} \int_0^\infty \lambda \left[ -\frac{Y_1^2}{\lambda^2} \right. \right. \\
 \left. \left. + \frac{2Z_1(\lambda)}{\lambda \rho_1} - \frac{(\lambda^2 + u_1^2 R_\perp(\lambda))}{\lambda^2} + (1 + R_\perp(\lambda)) \right] J_1(\lambda r) d\lambda \right\} \quad (2.3)
 \end{aligned}$$

Variables in these equations are defined as follows:

$\rho_1$ = resistivity of the first layer	$\lambda$ = dummy variable of integration,
$J_0$ = zeroth-order Bessel function,	$ds$ = source dipole length,
$J_1$ = first-order Bessel function,	$Y_m$ = $n^{\text{th}}$ layer wave number $(j\mu_n \omega \sigma_n)^{1/2}$ and
$I$ = current into source dipole,	$u_m^2 = \lambda^2 + Y_m^2$

$R_\perp(\lambda)$  and  $Z_\perp(\lambda)$  are correction factors which account for the effect of layering on the electromagnetic fields. The recursive expressions used for solving  $R_\perp(\lambda)$  and  $Z_\perp(\lambda)$  are written as follows:

$$R_\perp(\lambda) = \frac{N_0 - Y_1}{N_0 + Y_1} \quad (2.4) \quad \text{where } N_0 = \frac{\lambda}{j\mu_0 \omega}$$

$$Y_n = N_n \frac{Y_{n+1} + N_n \tanh(u_n h_n)}{N_n + Y_{n+1} \tanh(u_n h_n)}, \quad N_n = \frac{u_n}{j\mu_0 \omega}, \quad n=1, 2, \dots, N$$

$$\text{and } Y_N = \frac{u_N}{j\mu_0 \omega}$$

$$Z_\perp(\lambda) = K_n \frac{Z_{n+1} + K_n \tanh(u_n h_n)}{K_n + Z_{n+1} \tanh(u_n h_n)} \quad (2.5)$$

where  $K_n = u_n \rho_n$ ,  $n=1, 2, \dots, N$  and  $Z_N = u_N \rho_N$

The electric field components can be combined so that

$$E_c = \sin(2\theta) E_x - \cos(2\theta) E_y = \frac{I ds \rho_1 \sin(2\theta)}{4\pi} \int_0^\infty \gamma_1^2 (1 + R_1(\lambda)) + \frac{\lambda^2}{2} \left[ \frac{\gamma_1^2}{\lambda^2} + \frac{(\lambda^2 + u_1^2 R_1(\lambda))}{\lambda^2} - \frac{2 Z_1(\lambda)}{\lambda \rho_1} - (1 + R_1(\lambda)) \right] J_0(\lambda r) d\lambda \quad (2.6)$$

Combining  $E_x$  and  $E_y$  in this manner simplifies numerical evaluation and minimizes computational error.

To insure convergence of the integrals in equations (2.1), (2.2), (2.3), and (2.6) homogenous half-space expressions (assuming a resistivity of  $\rho_1$ ) are subtracted from the integrand. The closed-form expressions for a homogenous half-space are then added to the resulting integration. The final computational form for the field components then becomes:

$$E_x = -\frac{I ds \gamma_1^2 \rho_1}{4\pi} \int_0^\infty \left[ R_1(\lambda) - \frac{\lambda - u_1}{\lambda + u_1} \right] J_0(\lambda r) d\lambda + \frac{\rho_1 I ds}{4\pi} \left\{ -\cos^2 \theta \int_0^\infty \lambda^2 \left[ \frac{2}{\lambda} \left( \frac{Z_1(\lambda)}{\rho_1} - u_1 \right) - \frac{u_1^2}{\lambda^2} \left[ R_1(\lambda) - \frac{\lambda - u_1}{\lambda + u_1} \right] + R_1(\lambda) - \frac{\lambda - u_1}{\lambda + u_1} \right] J_0(\lambda r) d\lambda + \frac{\cos(2\theta)}{\pi} \int_0^\infty \lambda \left[ \frac{2}{\lambda} \left( \frac{Z_1(\lambda)}{\rho_1} - u_1 \right) + \left( 1 - \frac{u_1^2}{\lambda^2} \right) \left( R_1(\lambda) - \frac{\lambda - u_1}{\lambda + u_1} \right) \right] J_1(\lambda r) d\lambda + \frac{I ds \rho_1}{2\pi r^3} \left[ (3 \cos^2 \theta - 2) + (1 + \gamma r) e^{-\gamma_1 r} \right] \right\} \quad (2.7)$$

$$\begin{aligned}
E_y = & \frac{\rho_1 I ds}{4\pi} \left\{ -\sin(2\theta) \int_0^\infty \lambda^2 \left[ \frac{z_1(\lambda)}{\rho_1} - u_1 \right] \right. \\
& + \left( 1 - \frac{u_1^2}{\lambda^2} \right) \left( R_1(\lambda) - \frac{\lambda - u_1}{\lambda + u_1} \right) \left. \right] J_0(\lambda r) d\lambda + \frac{\sin(2\theta)}{\pi} \int_0^\infty \lambda \left[ \frac{z_1(\lambda)}{\rho_1} - u_1 \right] \\
& - \left( 1 - \frac{u_1^2}{\lambda^2} \right) \left( R_1(\lambda) - \frac{\lambda - u_1}{\lambda + u_1} \right) \left. \right] J_1(\lambda r) d\lambda + \frac{3\rho_1 I ds \sin\theta \cos\theta}{2\pi r^3} \quad (2.8)
\end{aligned}$$

$$\begin{aligned}
E_c = & \frac{\rho_1 I ds \sin(2\theta)}{2\pi} \left\{ \frac{1}{2} \int_0^\infty \left[ \frac{\gamma_1^2}{2} \left( \frac{\lambda - u_1}{\lambda + u_1} - R_1(\lambda) \right) + \lambda \left( u_1 - \frac{z_1(\lambda)}{\rho_1} \right) \right] J_0(\lambda r) d\lambda \right. \\
& + \frac{1}{r^3} \left[ -\frac{1}{2} + (1 + \gamma_1 r) e^{-\gamma_1 r} \right] \left. \right\} \quad (2.9)
\end{aligned}$$

$$\begin{aligned}
\text{and } H_z = & \frac{I ds \sin(\theta)}{2\pi} \left\{ \frac{1}{2} \int_0^\infty \lambda \left[ R_1(\lambda) - \frac{\lambda - u_1}{\lambda + u_1} \right] J_1(\lambda r) d\lambda \right. \\
& + \frac{1}{\gamma_1^2 r^4} \left[ 3 - (3 + 3\gamma_1 r + \gamma_1^2 r^2) e^{-\gamma_1 r} \right] \left. \right\} \quad (2.10)
\end{aligned}$$

To aid computations the following substitutions are made in equations (2.7), (2.8), (2.9), and (2.10):

$$\lambda = \frac{g}{\delta}, \quad K_m = \frac{\rho_1}{\rho_m}, \quad D_m = \frac{2h_m}{\delta}, \quad B = \frac{r}{\delta}$$

$$\text{and } u_m = \frac{V_m}{\delta} \quad \text{where } \delta^2 = \frac{2\rho_1}{\mu_0 \omega}, \quad \text{and } V_m^2 = g^2 + 2jK_m$$

Equations (2.7), (2.8), (2.9), and (2.10) are then written as:

$$\begin{aligned}
E_c = & \frac{I ds \sin(2\theta)}{4\pi \delta} \int_0^\infty \left\{ \gamma_1^2 \left( \frac{g V_1 (F_1 - 1)}{(g + V_1)(g + V_1 F_1)} + \frac{g V_1}{\delta^2} \right) J_0(gB) dg \right. \\
& + \frac{I ds \rho_1 \sin(2\theta)}{2\pi r^3} \left\{ -\frac{1}{2} + \left[ (1 + j)B + 1 \right] e^{-(1+j)B} \right\} \left. \right\} \quad (2.11)
\end{aligned}$$

$$\begin{aligned}
 E_x = & \frac{j \mu_0 \omega I ds}{4\pi} \left\{ \frac{\cos^2 \theta}{\delta} \int_0^\infty g \left[ i V_1(L, -1) + \frac{2 V_1(1-F_1)}{(g+V_1 F_1)(g+V_1)} \right] J_0(gB) dg \right. \\
 & - \frac{1}{\delta} \int_0^\infty \frac{2 g V_1(1-F_1)}{(g+V_1 F_1)(g+V_1)} J_0(gB) dg - \frac{\cos(2\theta)}{\pi} \int_0^\infty \left[ i V_1(L, -1) \right. \\
 & \left. \left. + \frac{2 V_1(1-F_1)}{(g+V_1 F_1)(g+V_1)} \right] J_1(gB) dg \right\} + \frac{I ds \rho_1}{2\pi \pi^3} \left[ (3 \cos^2 \theta - 2) + (1 + (1+j)B) e^{-(1+j)B} \right] \quad (2.12)
 \end{aligned}$$

$$\begin{aligned}
 E_y = & \frac{j \mu_0 \omega I ds}{4\pi} \left\{ - \frac{\sin(2\theta)}{\pi} \int_0^\infty \left[ j V_1(L, -1) + \frac{2 V_1(1-F_1)}{(g+V_1 F_1)(g+V_1)} \right] J_1(gB) dg \right. \\
 & + \frac{\sin(2\theta)}{2\delta} \int_0^\infty g \left[ j V_1(L, -1) + \frac{2 V_1(1-F_1)}{(g+V_1 F_1)(g+V_1)} \right] J_0(gB) dg \\
 & \left. + \frac{3 \rho_1 I ds}{2\pi \pi^3} \sin \theta \cos \theta \right\} \quad (2.13)
 \end{aligned}$$

$$\begin{aligned}
 H_z = & \frac{I ds}{2\pi \delta^2} \sin \theta \int_0^\infty \frac{g^2 V_1(1-F_1)}{(g+V_1 F_1)(g+V_1)} J_1(gB) dg \\
 & - \frac{j I ds \delta^2 \sin \theta}{4\pi \pi^4} \left\{ 3 - (3 + 3(1+j)B + 2jB^2) e^{-(1+j)B} \right\} \quad (2.14)
 \end{aligned}$$

$L_1$  and  $F_1$  are solved recursively, starting with the last layer and working upward, using

$$F_n = \frac{V_{n+1} F_{n+1} + V_n \left( \frac{1 - e^{-V_n D_n}}{1 + e^{-V_n D_n}} \right)}{V_n + V_{n+1} F_{n+1} \left( \frac{1 - e^{-V_n D_n}}{1 + e^{-V_n D_n}} \right)}, \text{ where } F_N = 1$$

$$\text{and } L_n = \frac{\rho_{n+1} V_{n+1} L_{n+1} + \rho_n V_n \left( \frac{1 - e^{-V_n D_n}}{1 + e^{-V_n D_n}} \right)}{\rho_n V_n + \rho_{n+1} V_{n+1} L_{n+1} \left( \frac{1 - e^{-V_n D_n}}{1 + e^{-V_n D_n}} \right)} \text{ where } L_N = 1.$$

## COMPUTATION OF THE ELECTROMAGNETIC FIELD COMPONENTS

Numerical Procedure

Computation of the integrals in equations (2.11), (2.12), (2.13), and (2.14) is carried out using the linear-digital-filter approach explained by Ghosh (1971). This technique is superior to other methods (e.g. spline integration (Vanyan, 1967), Gaussian quadrature (Anderson, 1973)) in its computational speed and ease of application. Linear filter integration yields about four significant figures of accuracy (Anderson, 1973).

The basis for linear filter integration is the sampling theorem. This theorem states that a function  $f(y)$  can be reconstructed by replacing each of the sampled values ( $y=a$ ) by

$$f(y) = \frac{f(a) \sin\left(\frac{2\pi(y-a)}{2\Delta y}\right)}{2\pi\left(\frac{y-a}{2\Delta y}\right)} \quad (3.1)$$

The function  $f(y)$  is equal to the infinite sum of the sinc-function weighted by  $f(a)$  ( $f(y) = \lim_{n \rightarrow \infty} \sum_{i=1}^n f(a_i) \text{sinc}(y)$ ).

The error involved in using this approximation is optimized when  $y = \frac{\ln(10)}{10}$  (Koeferd and others, 1972). This sample interval is used throughout the remainder of this paper. Substituting  $x = \ln(B)$  and  $y = \ln\left(\frac{1}{9}\right)$  into equations (2.11), (2.12), (2.13), and (2.14) yields

$$E_c = \frac{I ds \rho_s \sin(2\theta) e^{-x}}{4\pi \delta} \int_0^{\infty} E(\rho_m, h_m, y) e^{x-y} J_0(e^{x-y}) dy$$

$$+ \frac{I ds \rho_s \sin(2\theta)}{2\pi \pi^3} \left[ -\frac{1}{2} + \left( (1+j)B + 1 \right) e^{-(1+j)B} \right]$$

$$\begin{aligned}
E_x = & \frac{j \mu_0 \omega I ds}{4\pi} e^{-x} \left\{ \frac{\cos^2 \theta}{\delta} \int_{-\infty}^{\infty} E_{x1}(\rho_m, h_m, y) e^{x-y} J_0(e^{x-y}) dy \right. \\
& - \frac{1}{\delta} \int_{-\infty}^{\infty} E_{x2}(\rho_m, h_m, y) e^{x-y} J_0(e^{x-y}) dy - \frac{\cos(2\theta)}{\kappa} \int_{-\infty}^{\infty} E_{x3}(\rho_m, h_m, y) e^{x-y} J_1(e^{x-y}) dy \left. \right\} \\
& + \frac{I ds \rho_i}{2\pi \kappa^3} \left\{ [3 \cos^2 \theta - 2] + [1 + (1+j)B] e^{-(1+j)B} \right\}, \quad (3.3)
\end{aligned}$$

$$\begin{aligned}
E_y = & \frac{j \mu_0 \omega I ds}{4\pi} e^{-x} \left\{ -\frac{\sin(2\theta)}{\kappa} \int_0^{\infty} E_{y1}(\rho_m, h_m, y) e^{x-y} J_1(e^{x-y}) dy \right. \\
& + \frac{\sin(2\theta)}{2\delta} \int_0^{\infty} E_{y2}(\rho_m, h_m, y) e^{x-y} J_0(e^{x-y}) dy \left. \right\} + \frac{3\rho_i I ds \sin \theta \cos \theta}{2\pi \kappa^3}, \quad (3.4)
\end{aligned}$$

$$\begin{aligned}
\text{and } H_z = & \frac{I ds \sin \theta}{2\pi \delta^2} e^{-x} \int_0^{\infty} H(\rho_m, h_m, y) e^{x-y} J_1(e^{x-y}) dy \\
& - \frac{j I ds \delta^2 \sin \theta}{4\pi \kappa^4} \left\{ 3 - (3 + 3(1+j)B + 2jB^2) e^{-(1+j)B} \right\}, \quad (3.5)
\end{aligned}$$

$$\text{where } E(\rho_m, h_m, y) = \gamma_1^2 \left( \frac{e^{-y} v_1 (e^{-y}) [F_1 - 1]}{(e^{-y} + v_1 F_1)(e^{-y} + v_1)} + \frac{e^{-y} v_1 (1 - L_1)}{\delta^2} \right),$$

$$H(\rho_m, h_m, y) = \frac{e^{-2y} v_1 (1 - F_1)}{(e^{-y} + v_1 F_1)(e^{-y} + v_1)}, \quad E_{x2}(\rho_m, h_m, y) = \frac{2e^{-y} v_1 (1 - F_1)}{(e^{-y} + v_1 F_1)(e^{-y} + v_1)},$$

$$E_{y1}(\rho_m, h_m, y) = \frac{j v_1 (L_1 - 1) + 2 v_1 (1 - F_1)}{(e^{-y} + v_1 F_1)(e^{-y} + v_1)},$$

$$E_{x1}(\rho_m, h_m, y) = e^{-y} E_{y1}(\rho_m, h_m, y), \quad E_{y2}(\rho_m, h_m, y) = E_{x1}(\rho_m, h_m, y),$$

$$E_{x3}(\rho_m, h_m, y) = E_{y1}(\rho_m, h_m, y).$$

Equations (3.2), (3.3), (3.4), and (3.5) are in the form of convolution integrals.

If we let  $E(\rho_m, h_m, y) =$  input function,  $e^{(x-y)} J_0(e^{(x-y)}) =$  filter function, and apply the sampling theorem so that  $E(\rho_m, h_m, y) = E(a_i) \text{sinc}(y)$ , we obtain the relation

$$E(\rho_m, h_m, y) = \frac{I ds \rho_m \sin(2\theta)}{4\pi \delta} \sum_{i=1}^m E(a_i) \int_0^{\infty} \text{sinc}(y) e^{x-y} J_0(e^{x-y}) dy \quad (3.6)$$

Equation (3.6) can now be put into the following form:

$$E_c = \frac{I ds \rho_m \sin(2\theta)}{4\pi \delta} e^{-x} \sum_{i=1}^m E(a_i) c(y_i) \quad (3.7)$$

where  $c(y_i) = \text{sinc}(x) e^x J_0(x)$  and  $a_i = e^{-(x-y_i)}$ .  $c(y_i)$  is called the sinc-response of the filter function. The sinc-response is computed and stored as data in the computer program which calculates the field components.

$$\text{Using the identity } \int_0^{\infty} e^{-ay} J_0(yB) dy = \frac{1}{\sqrt{a^2 + B^2}} \quad (3.8)$$

and substituting  $x = \ln(B)$  and  $y = \ln(\frac{1}{B})$  the following integral is obtained

$$\int_0^{\infty} e^{-ae^{-y}} \left\{ e^{x-y} J_0(e^{x-y}) \right\} dy = \frac{e^x}{\sqrt{a^2 + e^{2x}}} \quad (3.9)$$

where "a" can be any positive constant. From equations (3.9) it follows that

$$\int_0^{\infty} [e^{-e^{-y}} - e^{-2e^{-y}}] \left\{ e^{x-y} J_0(e^{x-y}) \right\} dy = \frac{e^x}{\sqrt{1+e^{2x}}} - \frac{e^x}{\sqrt{4+e^{2x}}} \quad (3.10)$$

The sinc-response is now found using the expression

$$\text{sinc-response} = \mathcal{F}^{-1} \left[ \mathcal{F}(\text{sinc}) \mathcal{F} \left( \frac{\text{output function}}{\text{input function}} \right) \right] \quad (3.11)$$

where  $\mathcal{F}$  denotes the fourier transform operator and the inverse fourier transform operator,  $\mathcal{F}^{-1}$  denotes the input function equal to  $e^{-e^{-y}} - e^{-2e^{-y}}$ , and the output function

is equal to

$$\frac{e^x}{\sqrt{1+e^{2x}}} - \frac{e^x}{\sqrt{4+e^{2x}}}$$

Computations to obtain values for  $E_c$  are carried out in the following manner:

1. abscissa values "y<sub>i</sub>" are read into the computer and
2.  $E_c(\rho_m, h_m)$  is calculated using the relation

$$E_c(\rho_m, h_m) = \frac{I ds \rho_i \sin(2\theta) e^{-x}}{2\pi \delta} \sum_{i=1}^m E(a_i) c_i + \frac{I ds \rho_i \sin(2\theta)}{2\pi \pi^3} \left[ -\frac{1}{2} + ((1+j)B + 1) e^{-(1+j)B} \right] \quad (3.12)$$

The procedure for calculating  $H_z$ ,  $E_x$ , and  $E_y$  is similar to the method for calculating  $E_c$ . Putting  $H_z$  in the form of equation (3.7) we obtain

$$H_z(\rho_m, h_m) = \frac{I ds \sin \theta e^{-x}}{2\pi \delta^2} \sum_{k=1}^m H(b_k) d_k - j \frac{I ds \delta^2 \sin \theta}{4\pi \pi^4} \left\{ 3 - (3 + 3(1+j)B + 2jB^2) e^{-(1+j)B} \right\} \quad (3.13)$$

where  $d_k = \text{sinc}(x) * e^x J_1(x)$  and  $b_k = e^{-(x-Y_k)}$ . The sinc-response for  $H_z$  is calculated in the same manner as the  $E_c$  sinc-response except that the identity

$$\int_0^\infty q e^{-aq} J_1(qB) dq = \frac{B}{\sqrt{(a^2 + B^2)^{3/2}}}$$

is used instead of equation (3.8).

Using the procedure developed for  $E_c$  and  $H_z$  (zeroth and first order Hankel transforms) the final computational form for equations (3.3) and (3.4) becomes:

$$E_x(\rho_m, h_m) = j \mu_0 \omega I ds e^{-x} \left\{ \sum_{i=1}^m E_{x1}(a_i) c(y_i) - \frac{1}{\delta} \sum_{i=1}^m E_{x2}(a_i) c(y_i) - \frac{\cos(2\theta)}{\pi} \sum_{k=1}^m E_{x3}(b_k) d(y_k) \right\}$$

$$+ \frac{I ds \rho_1}{2 \pi r^3} \left\{ [3 \cos^2 \theta - 2] + [1 + (1+j)B] e^{-(1+j)B} \right\} \quad (3.14)$$

$$E_y(\rho_m, h_m) = j \frac{\mu_0 \omega I ds}{4 \pi} \left\{ -\frac{\sin(2\theta)}{r} \sum_{k=1}^{\infty} E_{y_1}(b_k) d(y_k) \right. \\ \left. + \frac{\sin(2\theta)}{2 \delta} \sum_{i=1}^{\infty} E_{y_2}(a_i) C(y_i) \right\} + \frac{3 \rho_1 I ds}{2 \pi r^3} \sin \theta \cos \theta \quad (3.15)$$

Computer Program for Calculating the Electromagnetic Field Components

A Fortran computer program to numerically calculate the equations for the electromagnetic field components is given in Appendix B. The program is called EMFWD. It computes  $E_x$ ,  $E_y$ ,  $E_z$ , and  $H_z$  using the numerical technique described in the previous section.

The results of this program are in agreement with results obtained by Anderson (1973). Figure 2 is a comparison of EMFWD results and results obtained by Anderson.

parameter set		EMFWD results	Anderson's results
A	$E_x$	$-.8296 \times 10^{-13} - j.1632 \times 10^{-12}$	$-.8291 \times 10^{-13} - j.1633 \times 10^{-12}$
	$E_y$	$.1310 \times 10^{-12} + j.2548 \times 10^{-12}$	$.1312 \times 10^{-12} + j.2548 \times 10^{-12}$
	$H_z$	$.1075 \times 10^{-11} - j.5272 \times 10^{-12}$	$.1075 \times 10^{-11} - j.5266 \times 10^{-12}$
B	$E_x$	$-.9804 \times 10^{-10} - j.1987 \times 10^{-10}$	$-.9811 \times 10^{-10} - j.1988 \times 10^{-10}$
	$E_y$	$.1570 \times 10^{-10} + j.3135 \times 10^{-10}$	$.1570 \times 10^{-10} + j.3134 \times 10^{-10}$
	$H_z$	$.6397 \times 10^{-09} - j.3018 \times 10^{-09}$	$.6398 \times 10^{-09} - j.3019 \times 10^{-09}$
C	$E_x$	$-.3190 \times 10^{-12} - j.5761 \times 10^{-12}$	$-.3190 \times 10^{-12} - j.5764 \times 10^{-12}$
	$E_y$	$.4999 \times 10^{-12} + j.9027 \times 10^{-12}$	$.5002 \times 10^{-12} + j.9026 \times 10^{-12}$
	$H_z$	$.7590 \times 10^{-12} - j.4212 \times 10^{-12}$	$.7591 \times 10^{-12} - j.4210 \times 10^{-12}$

parameter values:

set A;  $r=50000$ , frequency=1 hertz,  $\rho_1=100\Omega\cdot m$ ,  $\rho_2=1\Omega\cdot m$ ,  $h_1=1000$

set B;  $r=10000$ , frequency=1 hertz,  $\rho_1=100\Omega\cdot m$ ,  $\rho_2=1\Omega\cdot m$ ,  $h_1=1000$

set C;  $r=50000$ , frequency=1 hertz,  $\rho_1=100\Omega\cdot m$ ,  $\rho_2=1\Omega\cdot m$ ,  $h_1=1000$

for all sets  $I da=12.566371$

$r$  and  $h_1$  are in meters

Table 1 : Comparison of EMFWD results with values obtained by Anderson (personal communication, 1973).

Figures 2, 3, and 4 are examples of computational

results for the electric field components. Examples for the magnetic field component are given in Appendix A.

The theoretical sounding curves in figures 2 and 3 approach unity at high frequencies. In figure 2 the asymptotic behavior appears at a relatively low frequency. This effect is caused by the conductive first layer. Comparison of figures 3 and 4 indicates that the asymptotic behavior in figure 3 is caused by a thick first layer. Checks such as these establish the accuracy of EMFWD results and the validity of the numerical procedure which it utilizes.

Computation time for all complex electromagnetic field components is approximately 1.2 seconds per sample frequency (using the Colorado School of Mines DEC PDP-10 computer). EMFWD results are accurate to between 3 and 4 significant figures.

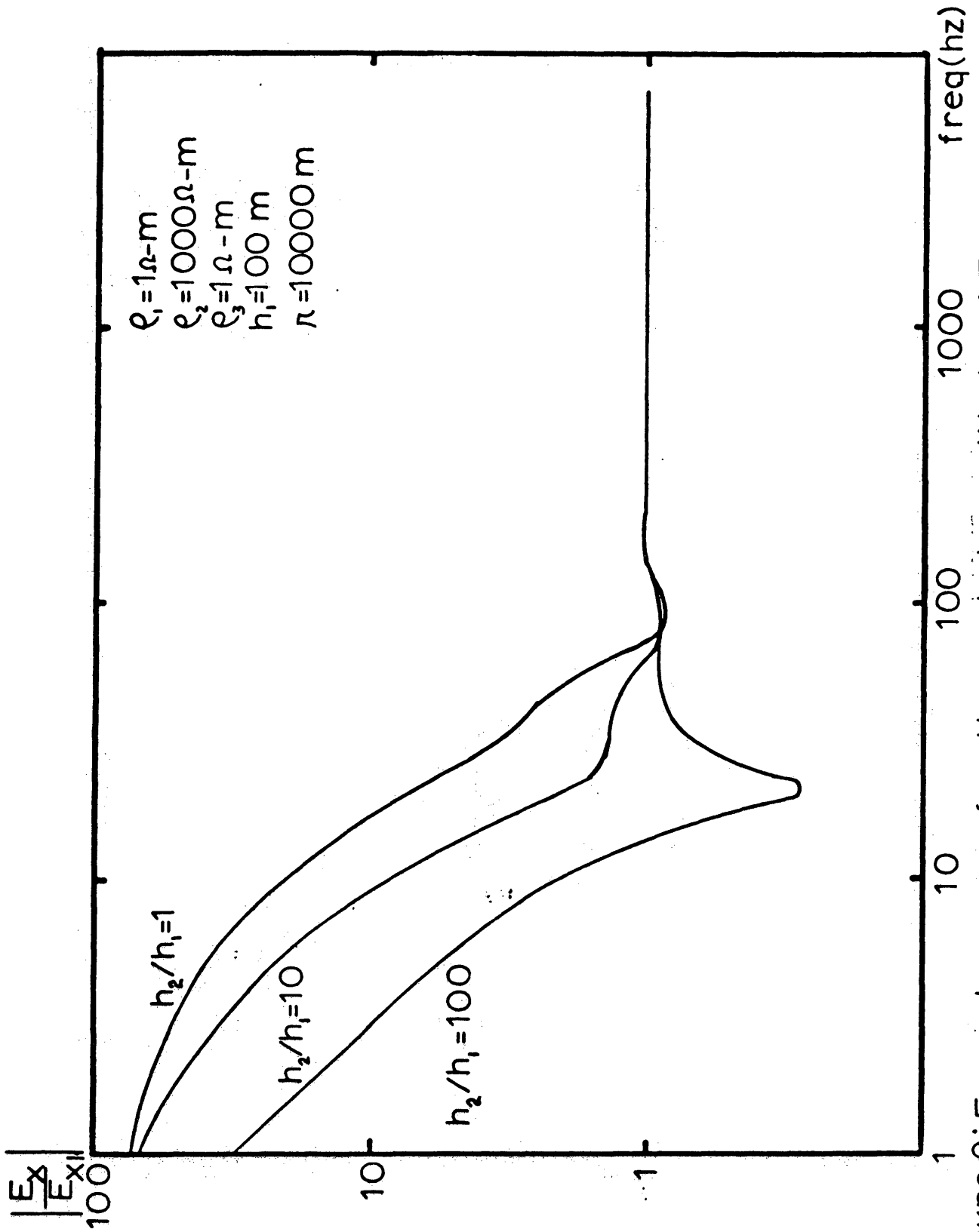


Figure 2: Example curve for the scaled amplitude of  $E_x$

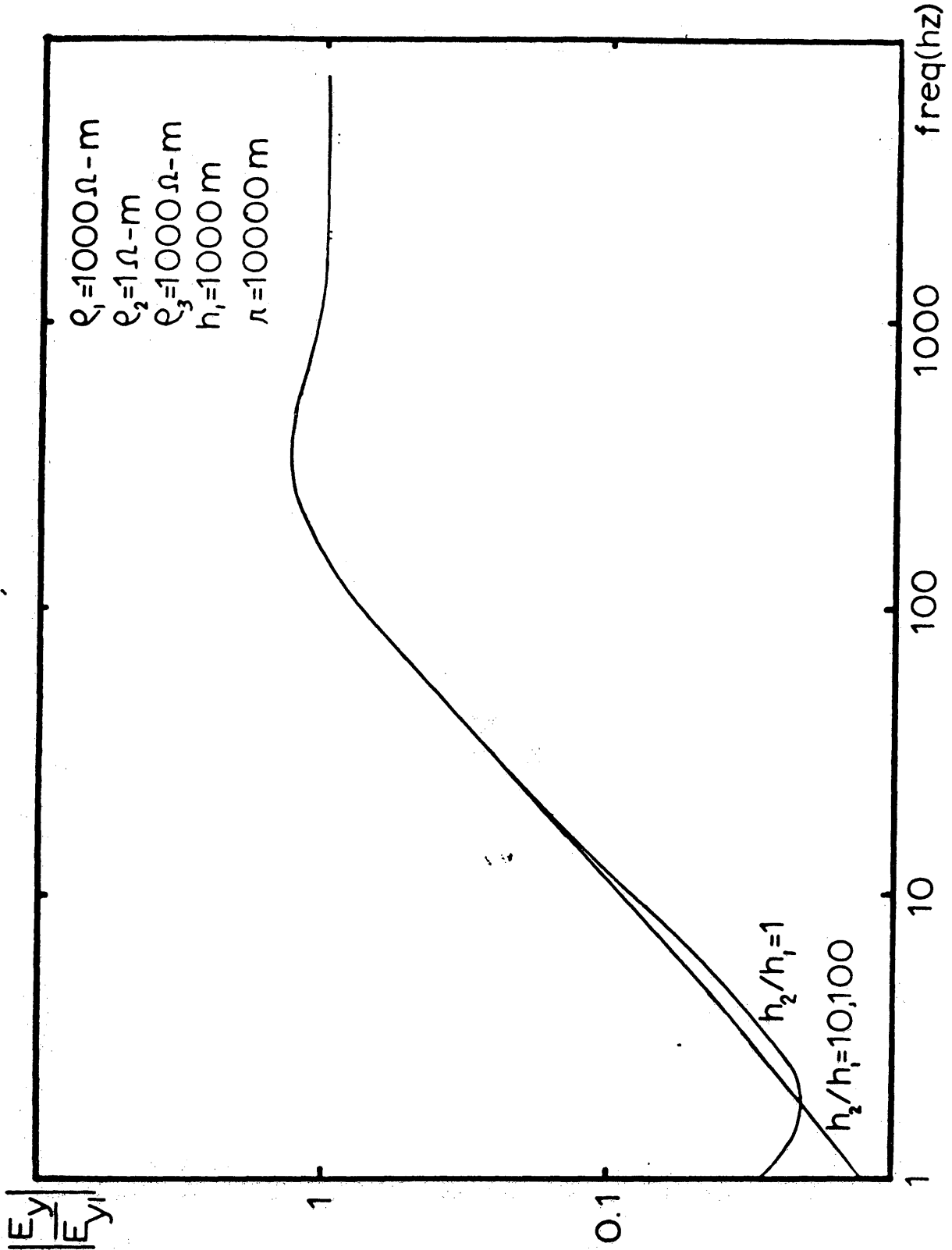


Figure 3: Example curve for the scaled amplitude of  $E_y$  ( $h_1 = 10000$ ).

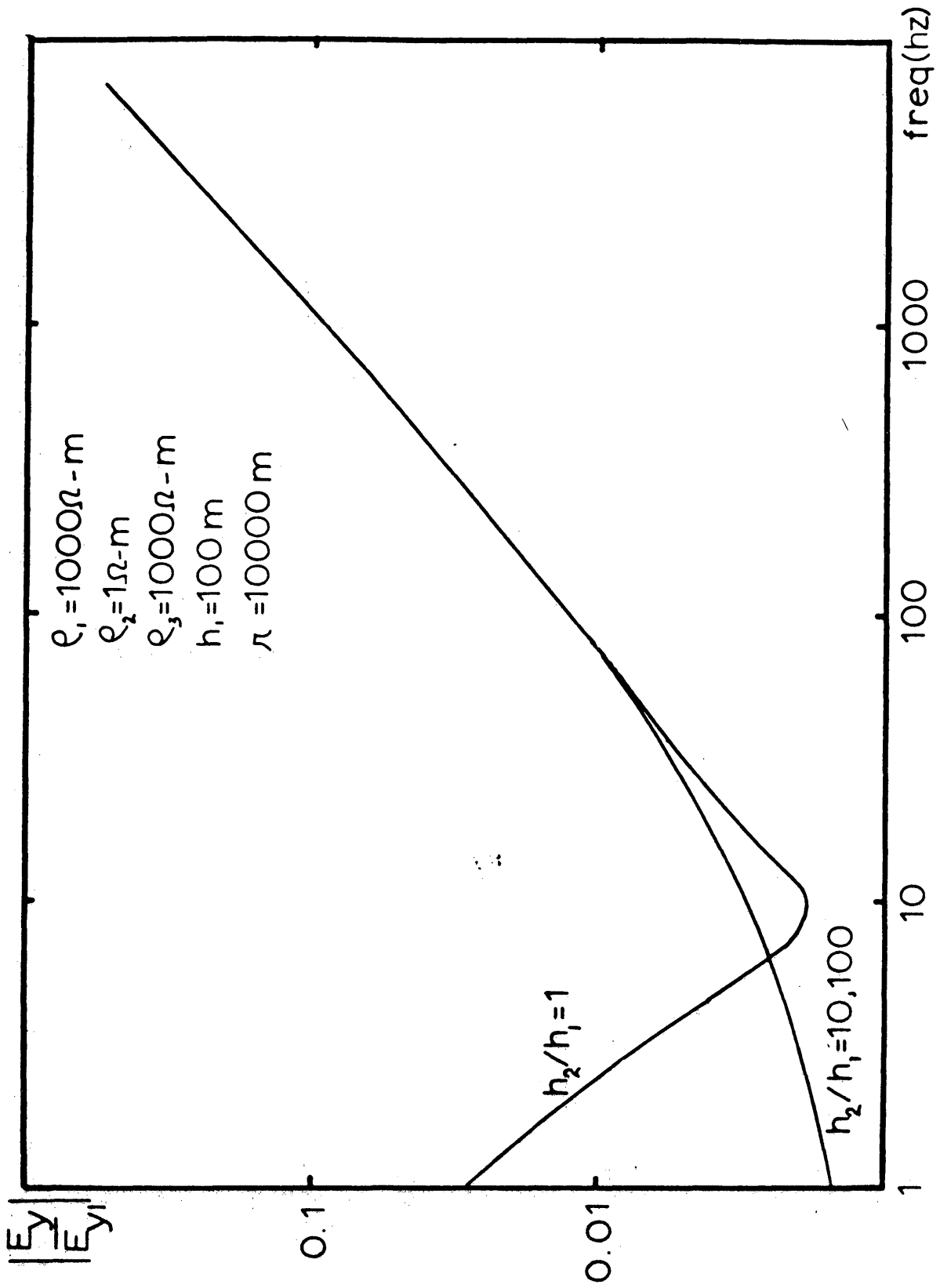


Figure 4: Example curve for the scaled amplitude of  $E_y$  ( $h_1=100$ ).

## LEAST-SQUARES CURVE FITTING

Theory of Least-Squares Curve Fitting

The objective of least squares curve fitting is to provide the best possible model approximation ( $\hat{Y}_i$ ) to the observed data ( $Y_i$ ). This is done by adjusting the parameters ( $P_j$ ) until  $\phi = \sum_{i=1}^m [Y_i - \hat{Y}_i]^2$  is minimized.  $\hat{Y}_i$  is the independent variable which is a function of  $x_i$  ( $i=1,2,\dots,n$ ) and  $P_j$  ( $j=1,2,\dots,k$ ). For the electromagnetic case considered in this paper  $\hat{Y}_i$  represents the theoretical field components and  $Y_i$  represents the observed field components at  $x_1, x_2, \dots, x_n$  frequencies. The parameters " $P_j$ " represent the layer thicknesses and layer resistivities.

The Newton-Gauss least-squares method approximates the field components using a Taylor series so that

$$\hat{Y}(x_i, \hat{P} + \delta_t) = f(x_i, \hat{P}) + \sum_{j=1}^k \left( \frac{\partial f_i}{\partial P_j} \right) \delta_{tj}, \quad (4.1)$$

where  $\hat{P}$  = least-squares estimate of  $P$ ,  $f(x_i, P) = \hat{Y}_i$  and  $\delta_t$  = correction vector to  $\hat{P}$ .

The least-squares criterion now becomes

$$\hat{\phi} = \sum_{i=1}^m [Y_i - f(x_i, \hat{P})] \sum_{j=1}^k \left( \frac{\partial f_i}{\partial P_j} \right) \delta_{tj} \quad (4.2)$$

$\delta_t$  is solved by setting  $\frac{\partial \hat{\phi}}{\partial \delta_j} = 0$  for all "j" values.

In matrix form the solution to the problem becomes

$$A \delta_t = \bar{q} \quad (4.3)$$

where  $A^{[k \times k]} = P^T P$ ,

$$P^{[m \times k]} = \frac{\partial f_i}{\partial P_j} \quad i = 1, 2, \dots, m; \quad j = 1, 2, \dots, k$$

$$\text{and } \mathbf{g}^{[k+1]} = \sum_{i=1}^n \left[ (Y_i - f_i) \frac{\partial f_i}{\partial P_j} \right] \quad (4.4)$$

In order to insure convergence with the Newton-Gauss method the first guess must be linearly close to the proper solution. The first guess is linearly close to the solution when the linearized model (equation 4.1) is a good approximation to the actual model.

To circumvent this problem Marquardt (1963) uses scaled matrices in the normal equations and introduces an adjustment factor,  $\lambda^n$ , which controls the convergence. According to the theory which he develops there always exists a  $\lambda$  (at the r-th iteration) such that  $\phi^{n+1} < \phi^n$ .

Following Marquardt's algorithm, the least-squares normal equations at the r-th iteration become

$$(\bar{A}^{*n} + \lambda^n \bar{I}) \bar{\delta}^{*n} = \bar{g}^{*n} \quad (4.5)$$

$$\text{where } \bar{A}^* = (a_{jj}^*) = \frac{a_{jj'}}{\sqrt{a_{jj'}} \sqrt{a_{jj'}}},$$

$$\bar{g}^* = g_j^* = \frac{g_j}{\sqrt{a_{jj'}}},$$

$$\text{and } \delta_j = \frac{\delta_j^*}{\sqrt{a_{jj'}}$$

Equation 4.5 is solved for  $\bar{\delta}^{*n}$  and the new parameter values become  $\hat{P}^{n+1} = \hat{P}^n + \delta^n$ . (4.6)

These parameter values are then used to compute a new set of theoretical points  $\hat{Y}_i$ , which are compared to the observed data. The iterative procedure is stopped when  $\phi$  is considered small enough for the particular problem being considered.

Marquardt suggests the following procedure be used to find the proper value of  $\lambda$ .

Let  $\nu > 1$ .

Let  $\lambda^{(n-1)}$  denote the value of  $\lambda$  from the previous iteration. Initially let  $\lambda^{(0)} = 10^{-2}$  say.

Compute  $\phi(\lambda^{(n-1)})$  and  $\phi(\lambda^{(n-1)}/\nu)$ .

i. If  $\phi(\lambda^{(n-1)}/\nu) \leq \phi^{(n)}$ , let  $\lambda^n = \lambda^{(n-1)}/\nu$ .

ii. If  $\phi(\lambda^{(n-1)}/\nu) > \phi^{(n)}$ , and  $\phi(\lambda^{(n-1)}) \leq \phi^{(n)}$ , let  $\lambda^n = \lambda^{(n-1)}$ .

iii. If  $\phi(\lambda^{(n-1)}/\nu) > \phi^{(n)}$ , and  $\phi(\lambda^{(n-1)}) > \phi^{(n)}$ , increase  $\lambda$  by successive multiplication by  $\nu$  until for some smallest  $w$ ,  $\phi(\lambda^{(n-1)}\nu^w) \leq \phi^{(n)}$ . Let  $\lambda^n = \lambda^{(n-1)}\nu^w$ .

### Application of Least-Squares Curve Fitting

Marquardt's algorithm is applied to electromagnetic sounding data as follows:

1. The appropriate theoretical model is chosen. In this paper a horizontally stratified earth is assumed to approximate the geological situation represented by the electromagnetic sounding data.

2. A first guess of the parameters is made. The parameters for a layered earth are electrical resistivity and layer thickness.

3.  $\phi^a$  is computed.

4. Equation (4.5) is solved for  $\delta^{*a}$ .

5. New parameters  $\hat{\rho}^{a+1}$  are computed according to equation (4.6).

6. A new theoretical solution is calculated using the parameters  $\hat{\rho}^{a+1}$ .

7.  $\phi^{a+1}$  is computed. At this point the computations are stopped if  $\phi^{a+1}$  is considered small enough. If  $\phi^{a+1}$  is not small enough a new  $\lambda$  is computed and steps 3 through 7 are repeated.

The computer program to apply Marquardt's algorithm to geophysical data was originally written by Jorge Parra (1972).

A modified version of this program, called EMINT, is listed in Appendix B.

### First Guess

Marquardt's algorithm requires an initial "guess" of the parameters. Obtaining an accurate first guess is critical for insuring fast and accurate convergence of the algorithm. A first guess can be obtained by catalog curve matching, by empirical "rules-of-thumb", or by using available geologic information.

Using the theoretical curves generated by EMFWD the following generalizations can be made concerning two and three layer curves:

(1) For a resistive first layer both the two and three layer  $\left| \frac{H_3}{H_{3,1}} \right|$  curves increase positively with increasing frequency. This effect becomes negligible for large first layer thicknesses. Figures (A-2) and (A-11) illustrate this behavior.

(2) For a conductive first layer  $\left| \frac{H_3}{H_{3,1}} \right|$  is equal to one at high frequencies. Figure (A-3) is an example of this condition.

(3) The ratio  $\left| \frac{H_3}{H_{3,1}} \right|$  for two and three layer curves is generally less than one at low frequencies for a resistive first layer and greater than one at low frequencies for a conductive first layer. See figures (A-1) and (A-12).

(4) Amplitude curves for the scaled electric field component approaches unity at high frequencies when the first layer is conductive. A thick resistive first layer will give the same effect. See figures 3 and 4 .

(4) At low frequencies the amplitude of the scaled electric field curves are less than unity for a resistive

first layer and greater than unity when the first layer is conductive. Figures (2.3), (2.4), and (2.5) illustrate this.

These rules alone are not sufficient to obtain an accurate first guess. A first guess made using these rules should be checked either with a curve catalog or with a test curve generated by EMFWD (Appendix B).

Appendix A contains a variety of two and three layer scaled  $H_z$  sounding curves. This catalog is by no means complete but it can serve to aid in obtaining a first guess. A major advantage of using scaled  $H_z$  field components

$$\frac{H_z}{H_z(p)} = \frac{\frac{1}{2} \int_0^{\infty} H(p_m, h_m, y) e^{x-y} J_0(e^{x-y}) dy + \frac{1}{\gamma^2 \rho^4} [3 - (3 + 3(1+j)B + 2jB) e^{-(1+j)B}]}{\frac{1}{\gamma^2 \rho^4} [3 - (3 + 3(1+j)B + 2jB) e^{-(1+j)B}]} \quad (4.7)$$

can be seen in their lack of dependence on the direction angle  $\theta$ , the source current  $I$ , and the dipole length  $ds$ . This simplification reduces the number of catalog curves required to make a first guess. Scaled  $E_c$  components have this same property. The scaled  $E_c$  components can be written in the form

$$\frac{E_c}{E_c(p)} = \frac{\frac{1}{2s} \int_0^{\infty} E(p_m, h_m, y) e^{x-y} J_0(e^{x-y}) dy + \frac{\rho_1}{\rho^3} \left\{ -\frac{1}{2} + [(1+j)B + 1] e^{-(1+j)B} \right\}}{\frac{\rho_1}{\rho^3} \left\{ -\frac{1}{2} + [(1+j)B + 1] e^{-(1+j)B} \right\}} \quad (4.8)$$

### Least-Squares Examples

Five examples (figures 5, 6, 7, 8 and 9 are shown to illustrate the results of least-squares curve fitting. The two layer examples show a resistive first layer overlaying a conductive basement while the three layer examples illustrate a resistive bed between two conductive beds.

Figure 5 shows a two layer example for the amplitude of  $H_z$ .  $\rho_1$  is fixed leaving only two variable parameters. Although it is not necessary to fix any parameters convergence is enhanced if either  $\rho_1$  or  $h_1$  is fixed. In this example the least-squares procedure was stopped after three iterations.

Table 2 shows the importance of choosing a good initial value for  $\lambda$ . From this example it appears that a smaller  $\lambda$  will yield a better fit. This is only true in cases where the first guess is close to the actual solution. When  $\lambda = 0$  Marquardt's algorithm is equivalent to the Newton-Gauss procedure which will diverge for many cases where the parameters are non-linear.

$\lambda$	$\phi$	RMS error
0.001	$207.45 \times 10^{-11}$	$3.38 \times 10^{-11}$
0.010	$214.40 \times 10^{-11}$	$3.44 \times 10^{-11}$
0.100	$277.70 \times 10^{-11}$	$3.92 \times 10^{-11}$
1.000	$643.40 \times 10^{-11}$	$5.96 \times 10^{-11}$
10.000	$1071.30 \times 10^{-11}$	$7.70 \times 10^{-11}$

where  $\phi = \sum_{i=1}^{\text{\# of points}} (Y_i - \hat{Y}_i)^2$ , RMS error =  $\sqrt{\frac{\phi}{\text{number of sample points}}}$

Table 2 : Effect of variation of  $\lambda$  on the interpretational error.

Figures 6 and 7 show a nearly perfect fit for the imaginary part of  $H_z$  and the real part of  $E_y$ . This fit was obtained after six iterations. Table 3 indicates the variation of error with respect to the number of iterations applied to this case. It is impossible to estimate the number of iterations which will yield a good fit for any given case. The following procedure will help to optimize the number of iterations necessary

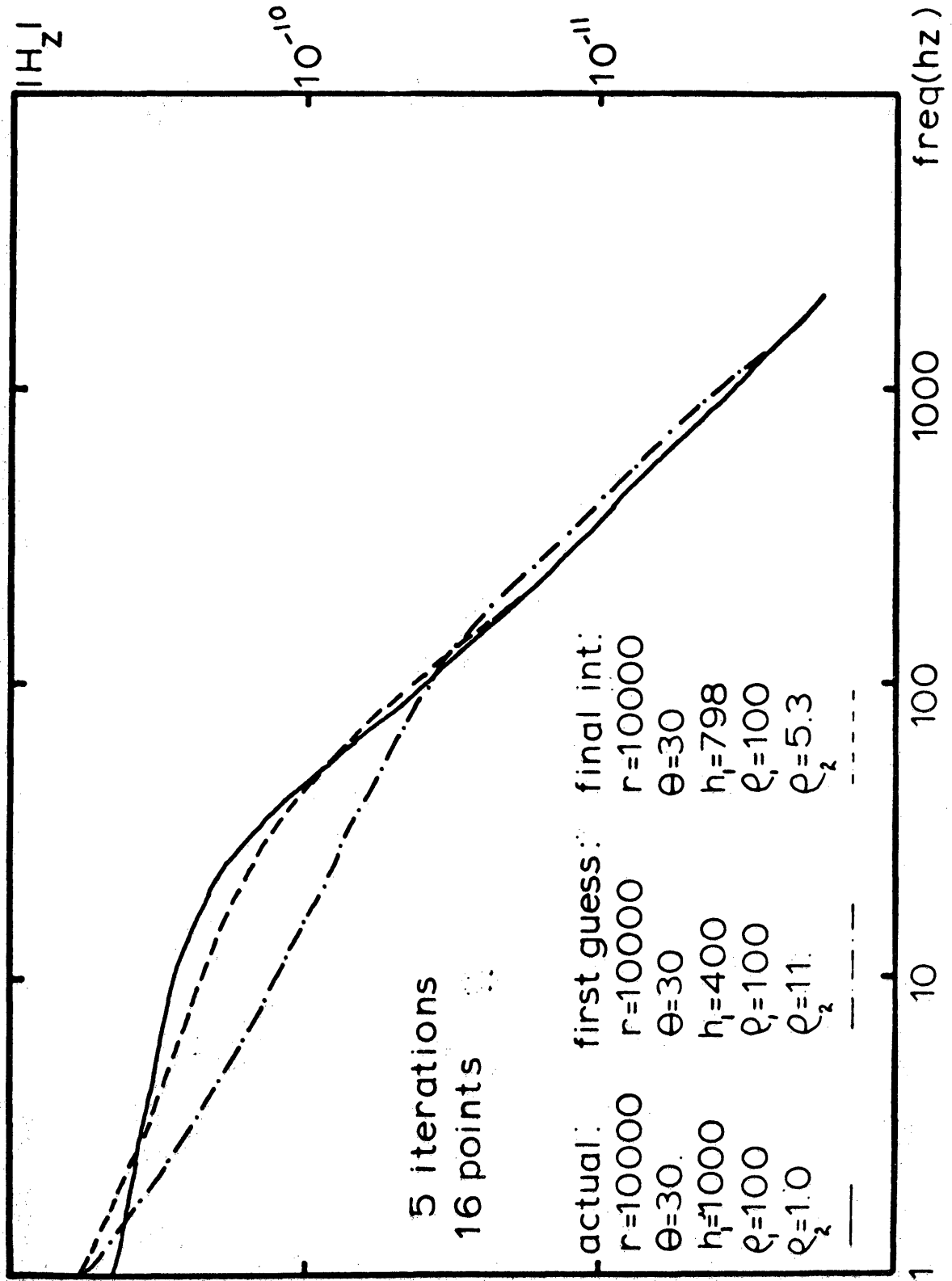


Figure 5: Two layer  $|H_2|$  least squares fit.

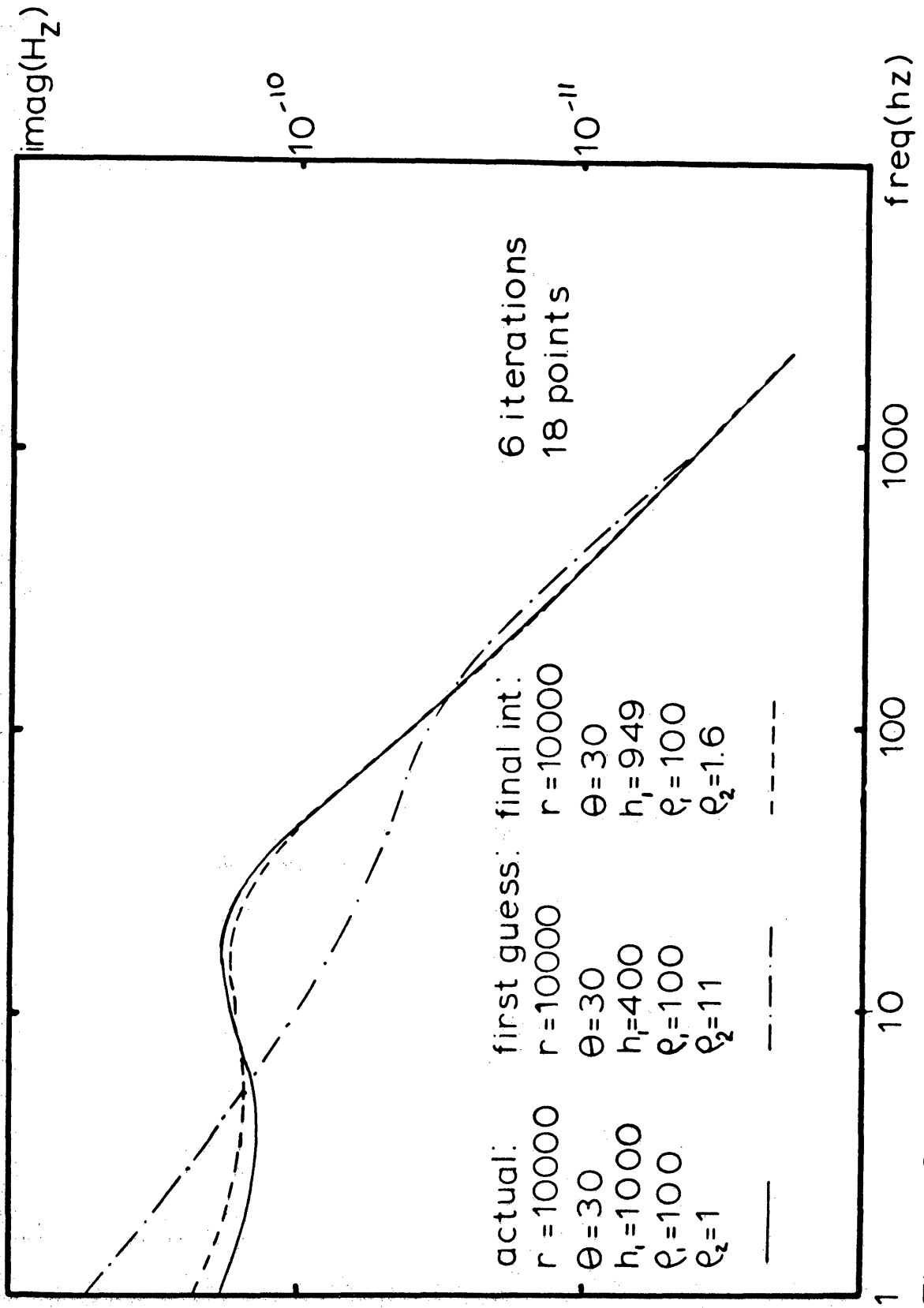


Figure 6: Two layer imaginary part of  $H_2$  least squares fit.

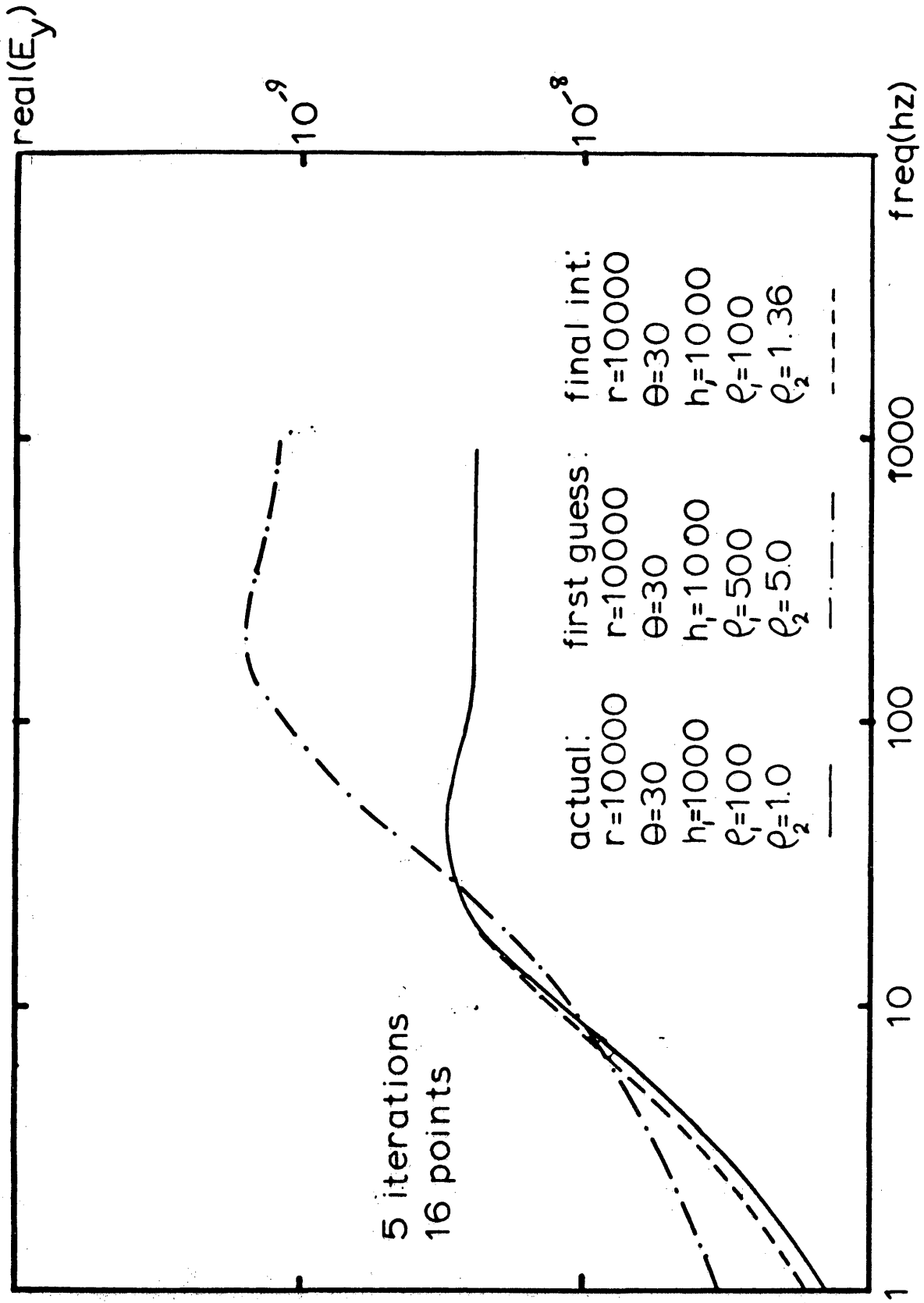


Figure 7 : Two layer real part of  $E_y$  least-squares curve fit.

to obtain a good fit.

Run EMINT through three iterations at a time.

i. If  $\phi_3 < \phi_2$  by more than ten-percent, continue the iterative procedure.

ii. If  $\phi_2 < \phi_1$  by more than ten-percent and  $\phi_3 < \phi_2$  by less than ten-percent check the closeness of fit between the observed data and the calculated values. Stop the iterative procedure if the fit is considered adequate. Change  $\lambda$  or the first guess if the fit is inadequate.

iii. If  $\phi_2 < \phi_1$  by less than ten-percent and  $\phi_3 < \phi_2$  by less than ten-percent change  $\lambda$  or the first guess and continue the iterative procedure.

<u>Iteration</u>	<u><math>\phi</math></u>	<u>RMS error</u>	<u>(ohm-m)</u>	<u>(meters)</u>
3	$306.3 \times 10^{-10}$	$10.1 \times 10^{-10}$	3.30	726.4
4	$182.7 \times 10^{-10}$	$2.12 \times 10^{-10}$	2.64	823.4
6	$49.1 \times 10^{-10}$	$2.9 \times 10^{-11}$	1.68	949.4
Actual parameter values			1.00	1000.0
First guess parameter values			11.00	400.0

Table 3 : Variation of converged parameters and error with the number of iterations.

Table 4 indicates how the accuracy of fit is affected by the number of sample points used for interpretation. Increasing the number of sample points does not necessarily improve the fit. As long as the curve is adequately defined, only about three sample points per decade of frequency are required for a good fit. This conclusion is based solely on theoretical results. More points are necessary when

the data contains geologic or other non-random noise.

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**Two Layers :  $Imag(E_y)$**

Actual parameters  $r=10000m., \theta=30^\circ, \rho_1=100\Omega\cdot m, h_1=100m., \rho_2=1.\Omega\cdot m$

First guess  $r=10000m., \theta=30^\circ, \rho_1=500\Omega\cdot m, h_1=100m., \rho_2=5.\Omega\cdot m$

minimum frequency= 1hertz maximum frequency= 1000hertz

all points evenly spaced

<u>Number of points</u>	<u>RMS error</u>	<u>Interpreted parameters</u>	
6	$1.06 \times 10^{-11}$	$\rho_1=100.02\Omega m$	$\rho_2=1.142 m$
8	$1.13 \times 10^{-11}$	$\rho_1=100.01\Omega m$	$\rho_2=1.145 m$
10	$1.07 \times 10^{-11}$	$\rho_1=100.01\Omega m$	$\rho_2=1.137 m$
12	$1.22 \times 10^{-11}$	$\rho_1=100.01\Omega m$	$\rho_2=1.148 m$

**Three Layers :  $|H_z|$**

Actual parameters  $r=10000m., \theta=30^\circ, \rho_1=1.\Omega\cdot m, \rho_2=1000.\Omega\cdot m,$   
 $\rho_3=1.\Omega\cdot m, h_1=100m., h_2=100m.$

First guess  $r=10000m., \theta=30^\circ, \rho_1=1.\Omega\cdot m, \rho_2=1300.\Omega\cdot m,$   
 $\rho_3=5.\Omega\cdot m, h_1=200m., h_2=300m.$

minimum frequency=1 hertz maximum frequency=6000 hertz

all points evenly spaced

<u>Number of points</u>	<u>RMS error</u>	<u>Interpreted parameters</u>	
7	$1.63 \times 10^{-12}$	$h_1=196.2 m$	$h_2=134.0 m$
		$\rho_2=1121.5 \Omega m$	$\rho_3=1.38 \Omega m$
13	$1.86 \times 10^{-12}$	$h_1=196.4 m$	$h_2=147.9 m$
		$\rho_2=1126.4 \Omega m$	$\rho_3=1.53 \Omega m$
20	$1.78 \times 10^{-12}$	$h_1=187.4 m$	$h_2=156.7 m$
		$\rho_2=1177.4 \Omega m$	$\rho_3=1.111 \Omega m$

**Table 4 :** The effect of varying the number of sample points on the interpretation error and interpreted parameter values.

---

Figures 8 and 9 show test results of applying Marquardt's least-squares algorithm to three layer cases. For the three layer case proper convergence can only occur when the first guess is reasonably close to the exact solution.

Computer time and cost for the least-squares interpretation program can be summarized as follows (cost and time per sample point):

(1) For  $H_z$  and  $E_c$  two layer cases each iteration takes approximately two seconds (\$1.20) while  $E_x$  and  $E_y$  takes four seconds (\$2.40) of computer time.

(2) For  $H_z$  and  $E_c$  three layer cases each iteration takes approximately three seconds (\$1.80) while  $E_x$  and  $E_y$  takes five seconds (\$3.00) of computer time.

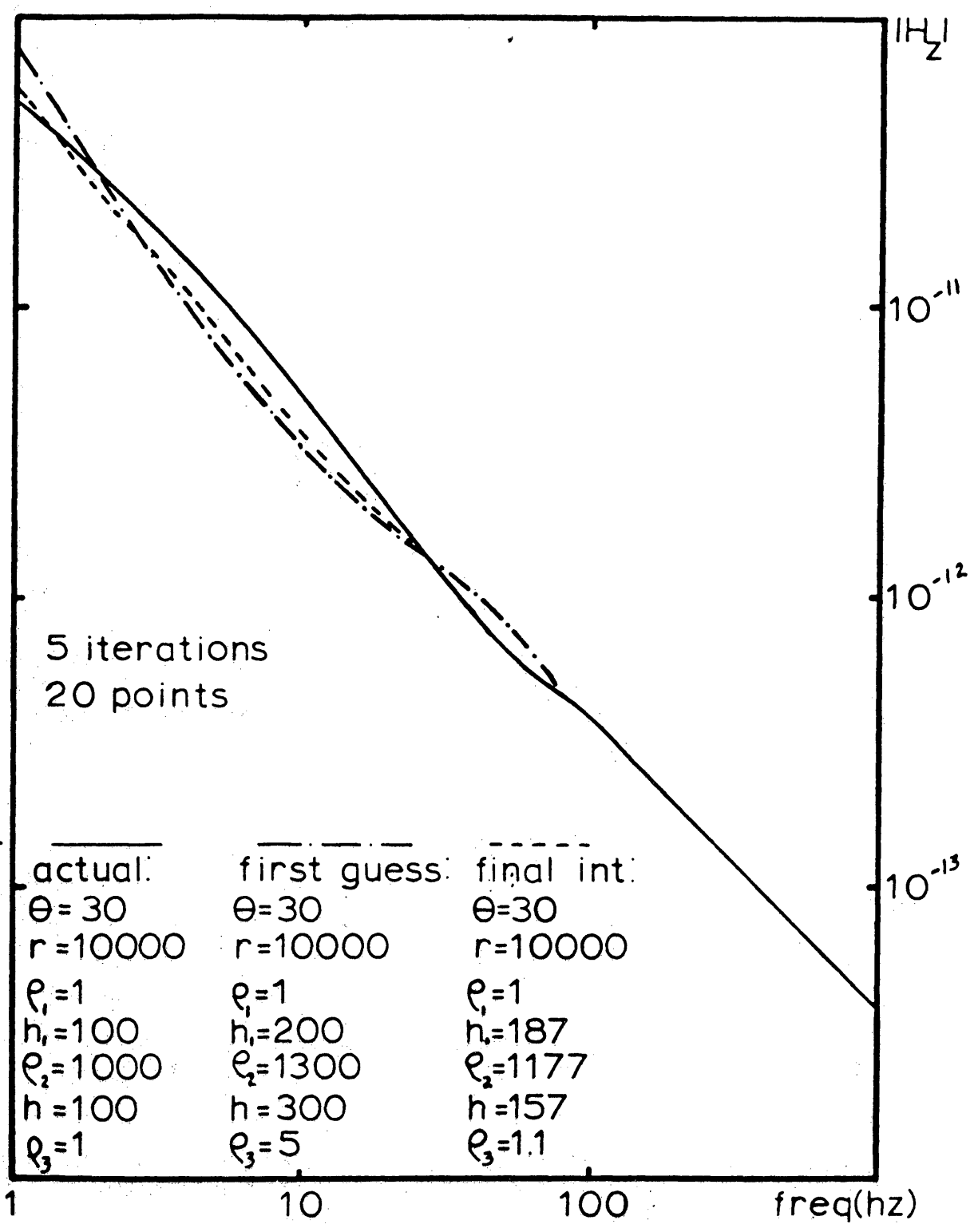


Figure 8: Three layer  $|H_2|$  least squares-curve fit.

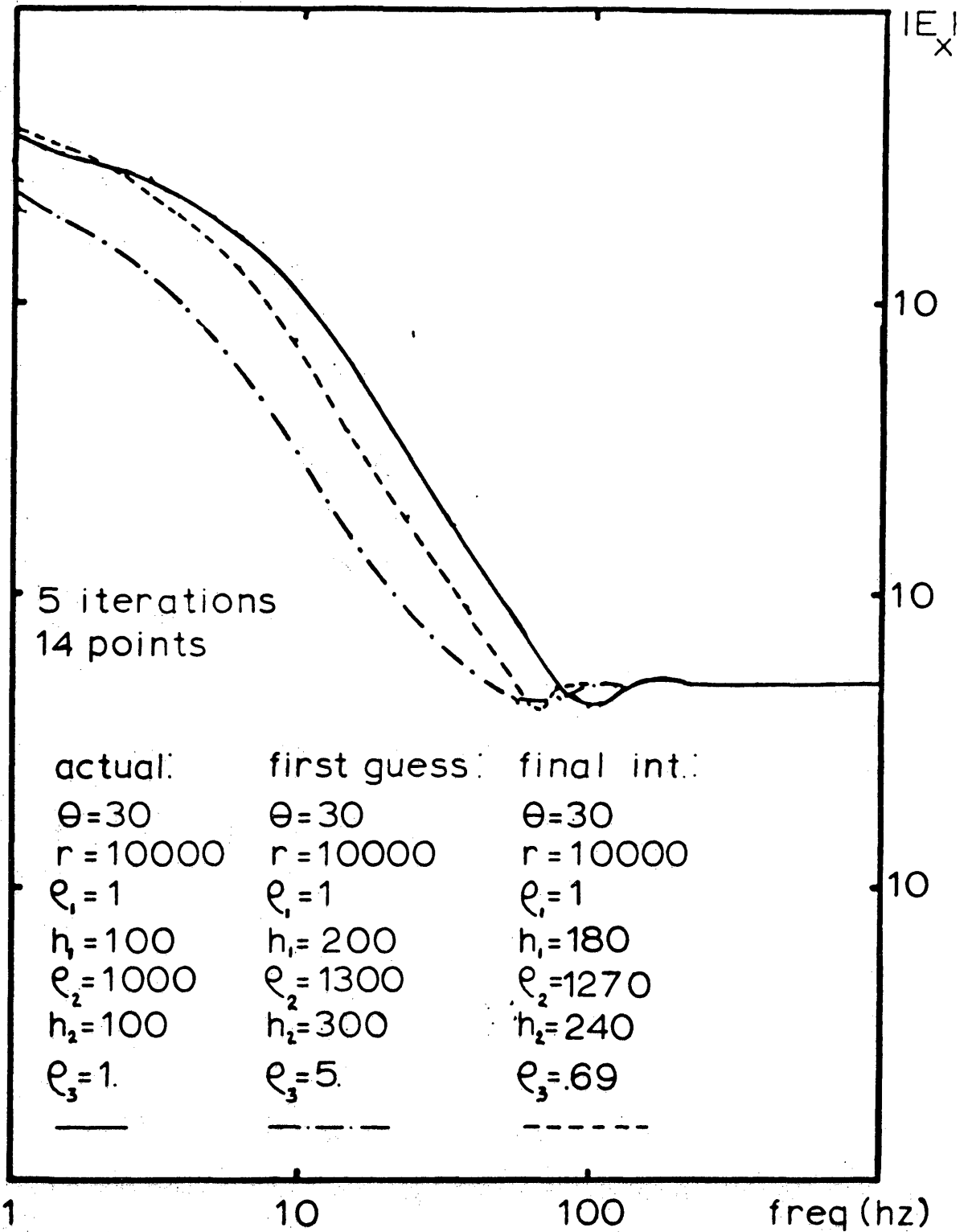


Figure 9: Three layer  $|E_x|$  least-squares curve fit.

## CONCLUSIONS

Marquardt's algorithm for least-squares curve fitting provides a good method for interpreting electromagnetic sounding data. The theoretical expressions for the electromagnetic field components can be evaluated numerically by making the proper substitution of variables in the integral expressions and applying principles of linear filter theory to the resulting integrals.

Test results of applying Marquardt's algorithm to electromagnetic sounding data can be summarized as follows:

(1) For the two layer case convergence of the layered earth parameters is generally fast and accurate. Convergence of the layered earth parameters ( resistivity and thickness) is accurate to within ten-percent. Convergence is enhanced by fixing either the first layer resistivity or the first layer thickness.

(2) Three layer curve fitting, with four or more parameters, requires a good first guess to insure convergence. Parameter convergence is generally slower for the three layer case than for the two layer case.

Scaled  $H_z$  field component curves (Appendix A ) can simplify catalog curve matching by eliminating the curves dependence on  $\theta$ ,  $I$ , and  $dc$ . Similarity, scaling  $E_c$  ( which is a mathematical combination of  $E_x$  and  $E_y$ ) simplifies catalog interpretation for the  $E_x$  and  $E_y$  components of the electric field. These modifications of the field components can enhance the ability of the interpreter to obtain a good first guess.

The overall least-squares computer program is too expensive to operate on a commercial basis. For  $H_z$  and  $E_c$  two layer cases, each iteration takes approximately two seconds per sample point, while for  $E_x$  and  $E_y$  each iteration consumes four seconds of computer time per sample point. Computation time can be reduced by calculating the partial derivatives of the layered-earth parameters in closed form.

Further computer time could be saved by using separate computer programs to individually fit each field component.

Regardless of how "automated" an interpretation scheme may be, meaningful results still require good judgment on the part of the interpreter. The geophysicist who uses the technique developed in this paper should not lose sight of the assumptions which are inherent in the method. The final interpretation of electromagnetic sounding data must be geologically meaningful.

APPENDIX A

Two and Three Layer Scaled  $H_z$  Amplitude Curves

The graphs in this section were computed using Program EMFWD. They may serve as a guide for a first guess interpretation.

Graph symbols:

$$\frac{H}{H_1} = \frac{\text{amplitude of } H_z \text{ (layered earth model)}}{\text{amplitude of } H_z \text{ (homogenous half-space)}}$$

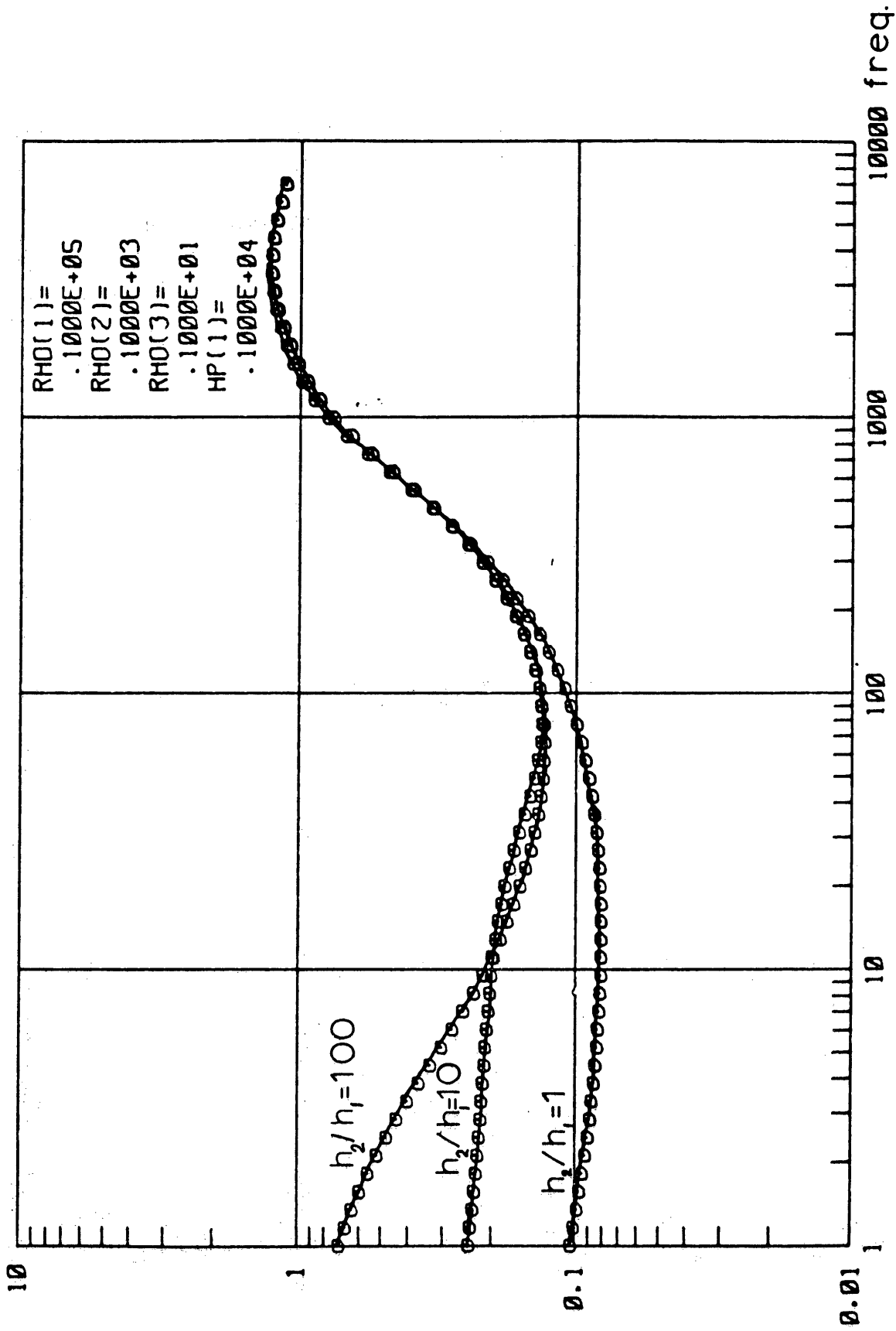
$\frac{H}{H_1}$  represents the ordinate values on each graph.

R = source receiver separation

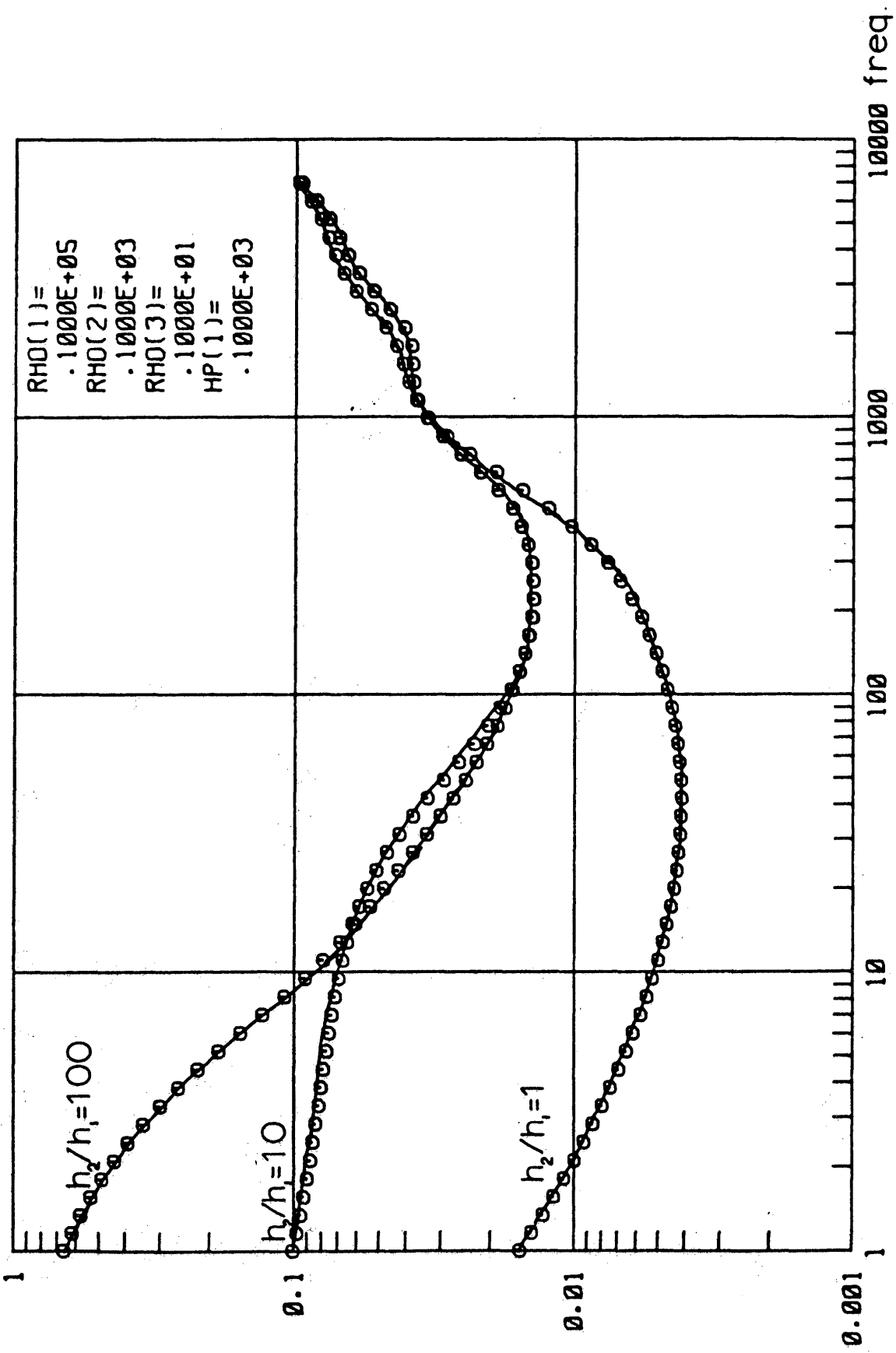
RHO(i) = resistivity of the i-th layer

HP(i) = thickness of the i-th layer

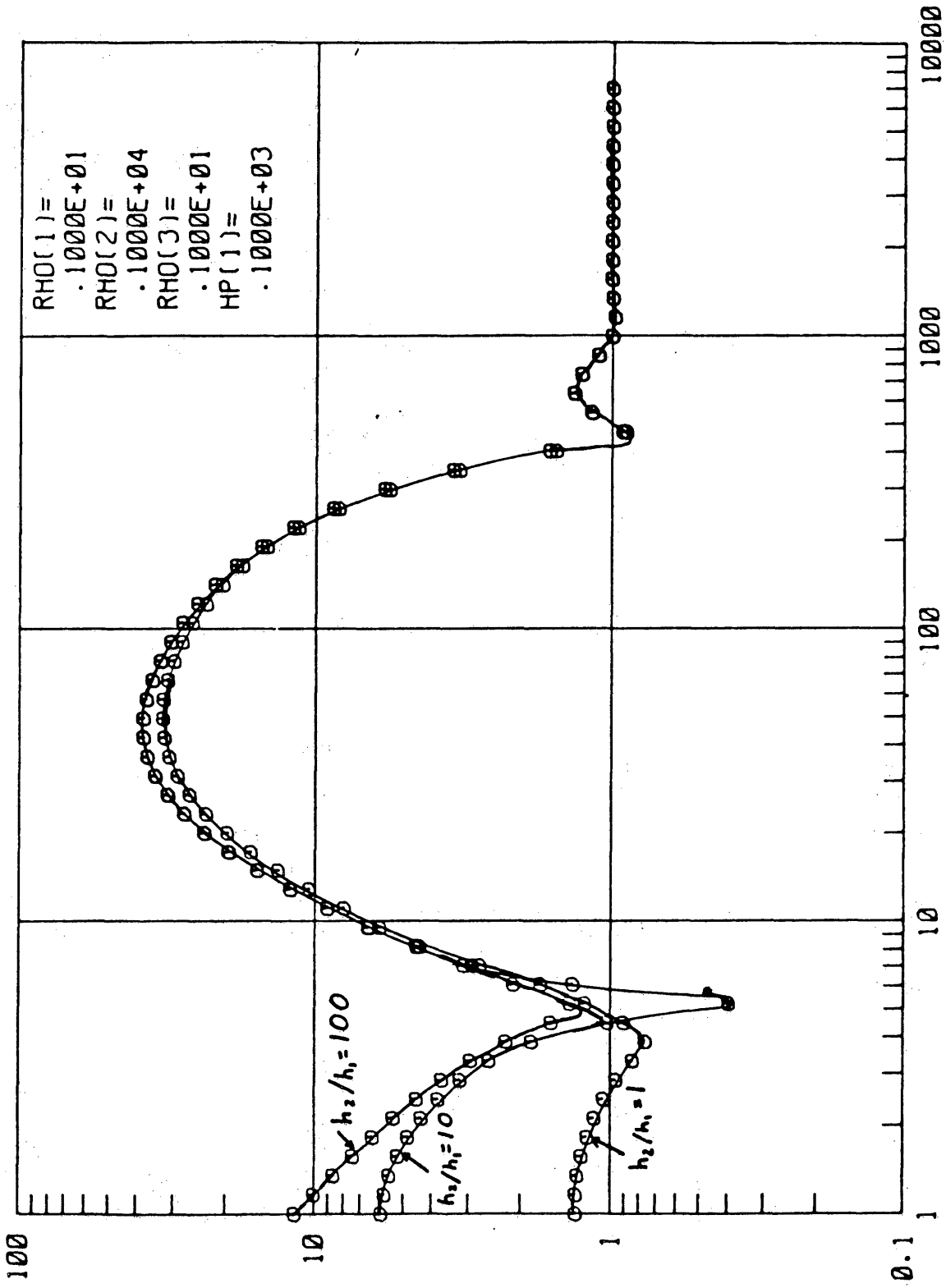
abscissa values = frequency (in hertz)



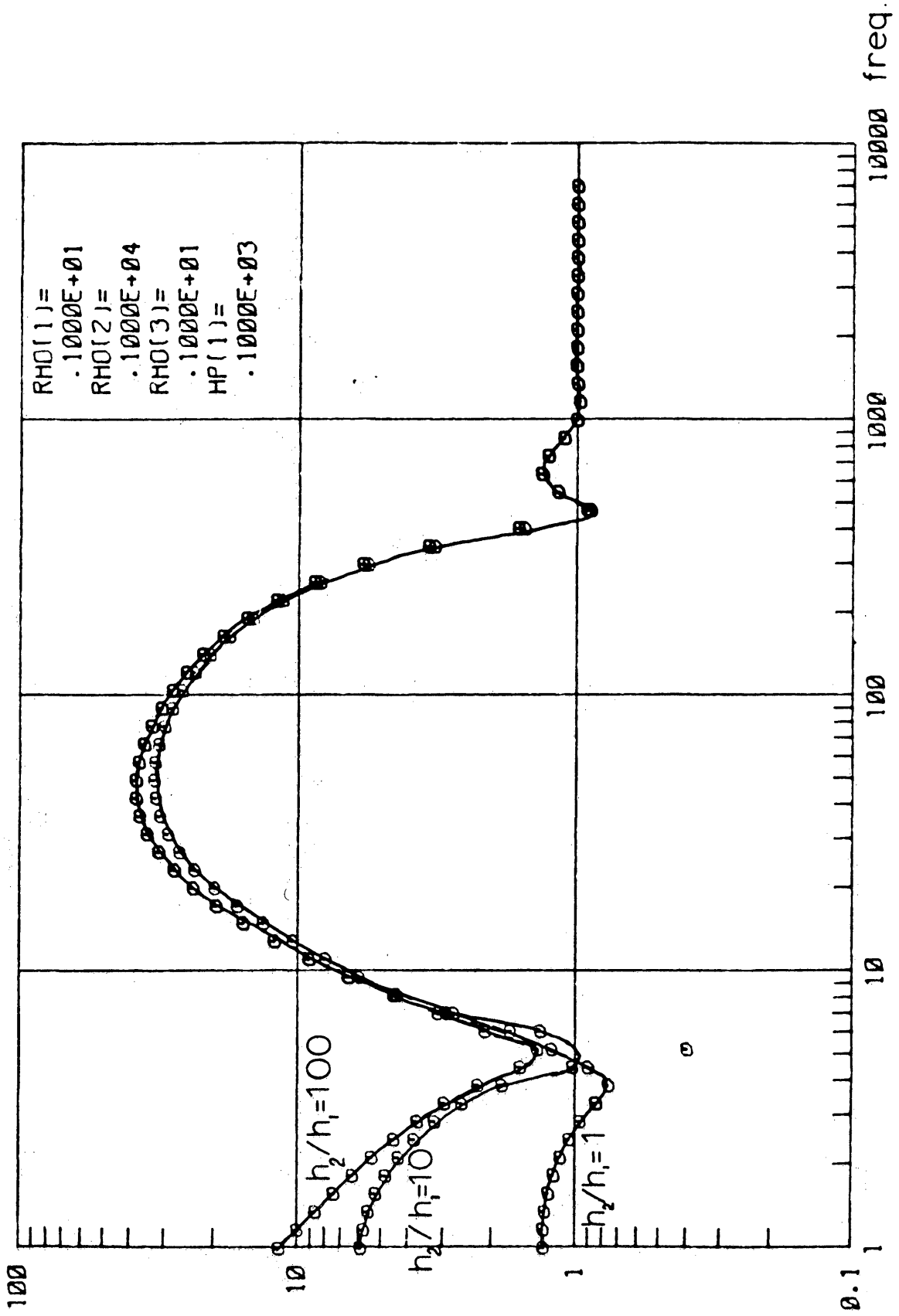
|H/H1| 3 LAYERS R = .100E+05



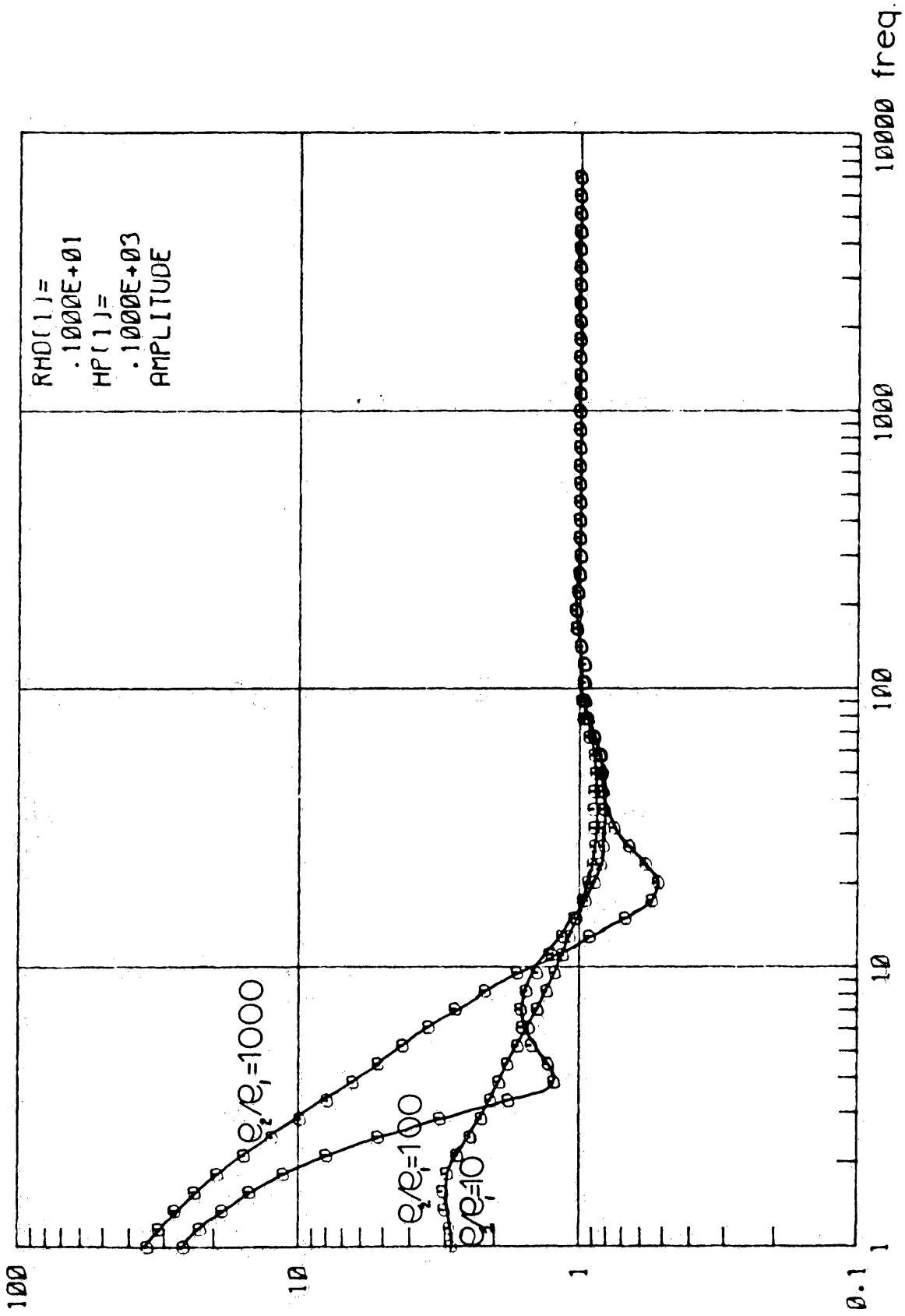
IH/H1I 3 LAYERS R = .100E+05



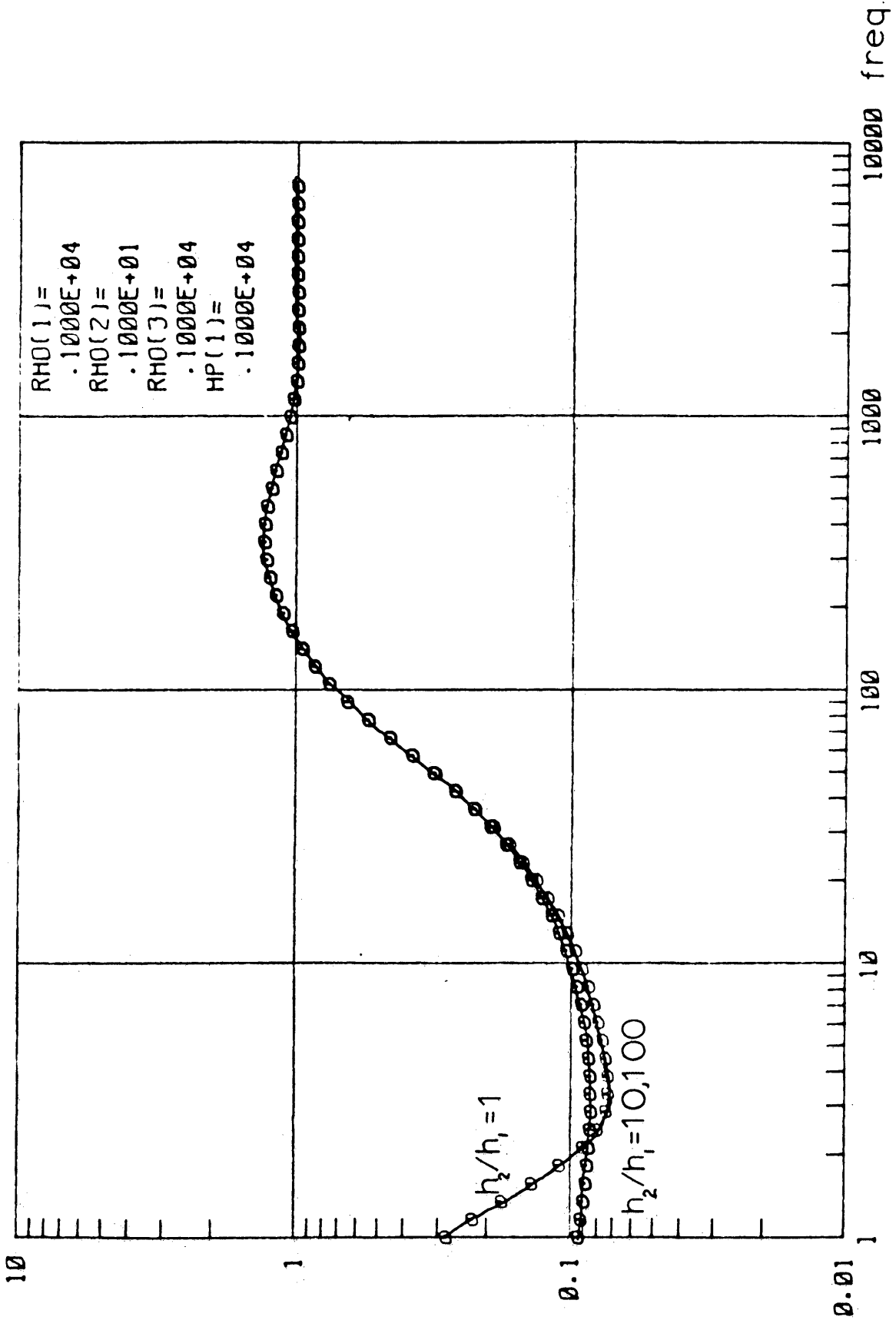
|H/H1| 3 LAYERS R = .500E+05



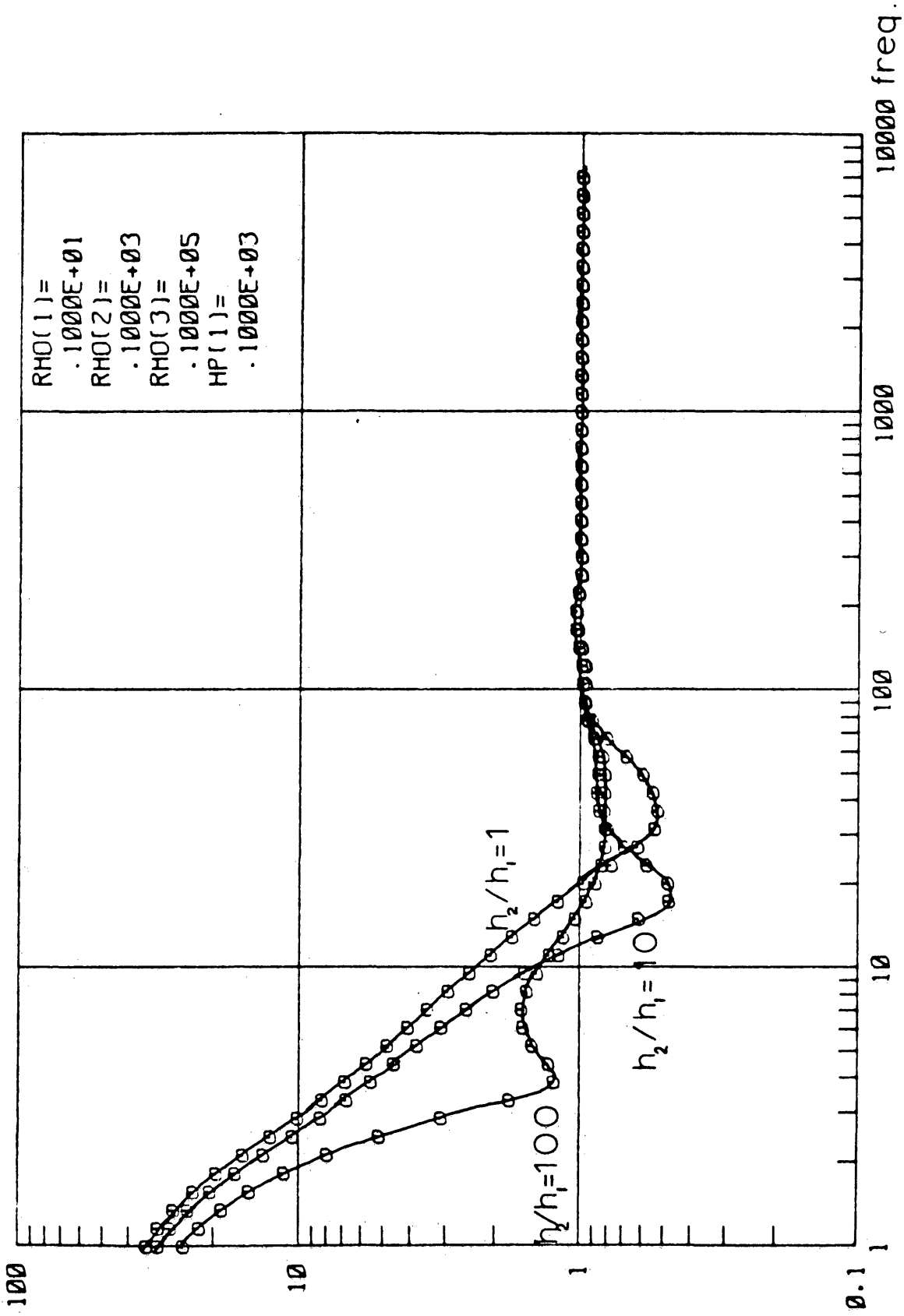
|H/H1| 3 LAYERS R = .500E+05



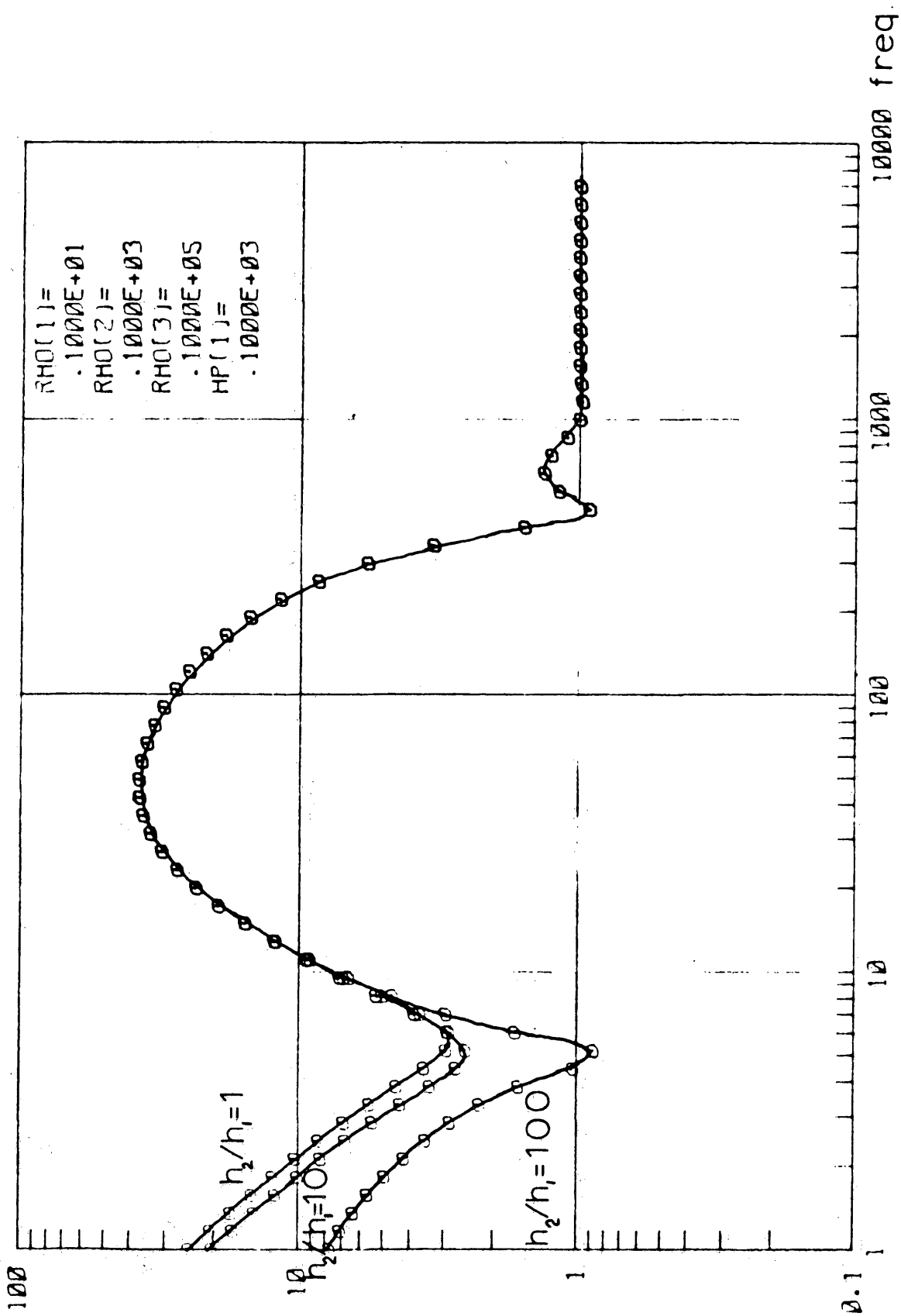
|H/H1| 2 LAYERS R = .100E+05



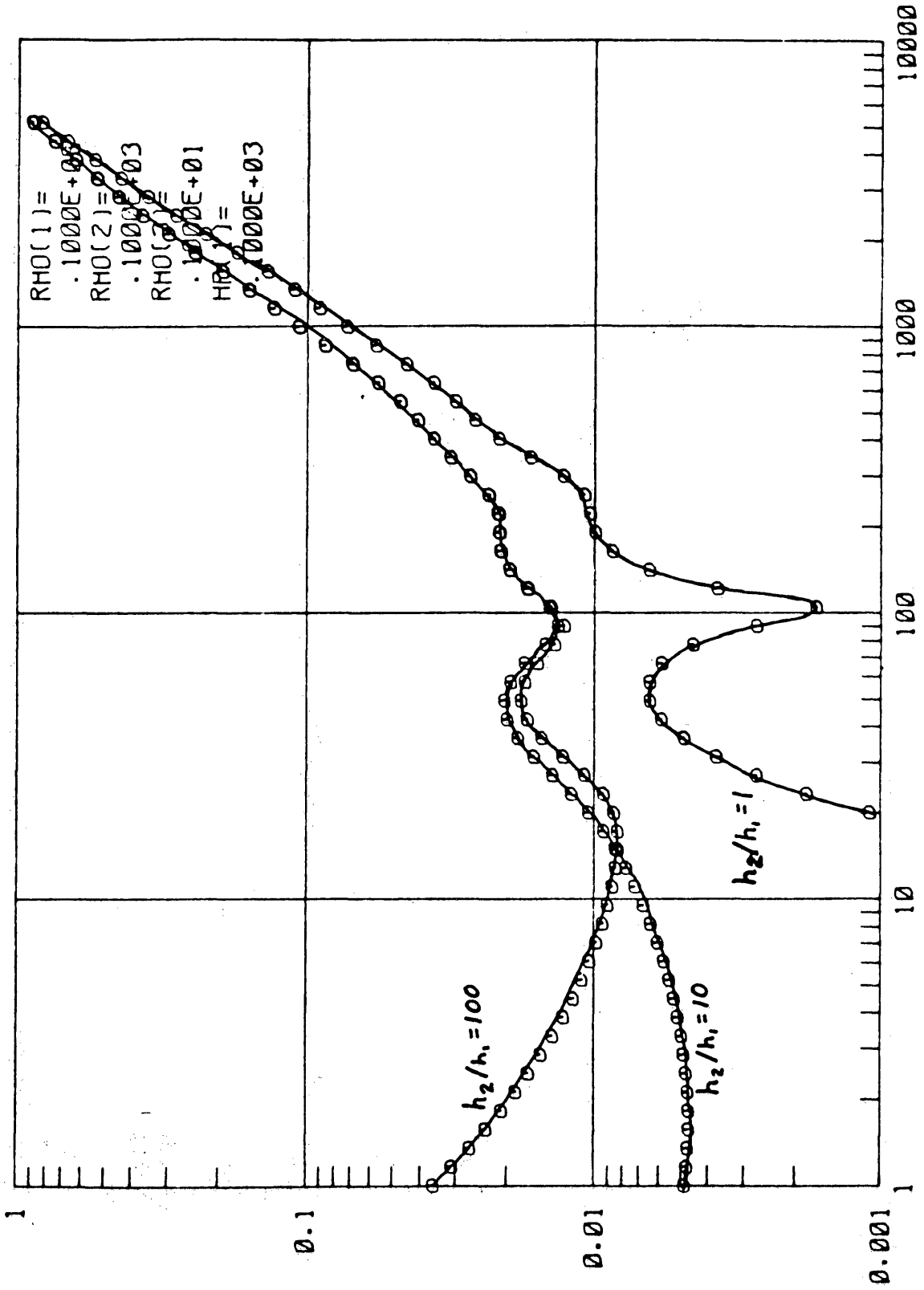
H/H11 3 LAYERS R = .100E+05



$h/h_1$  3 LAYERS  $R = .100E+05$

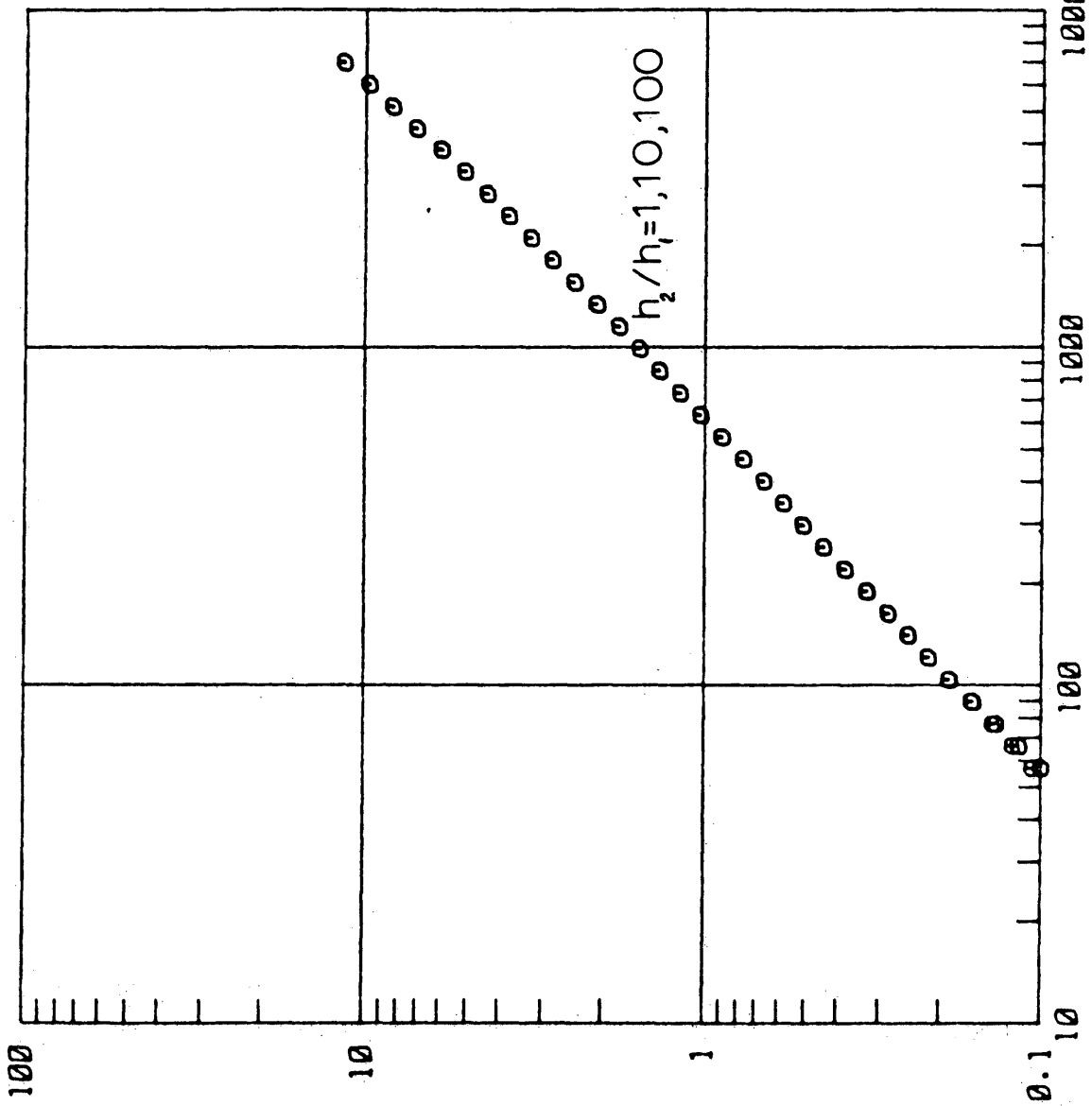


H/H1 3 LAYERS R = .500E+05

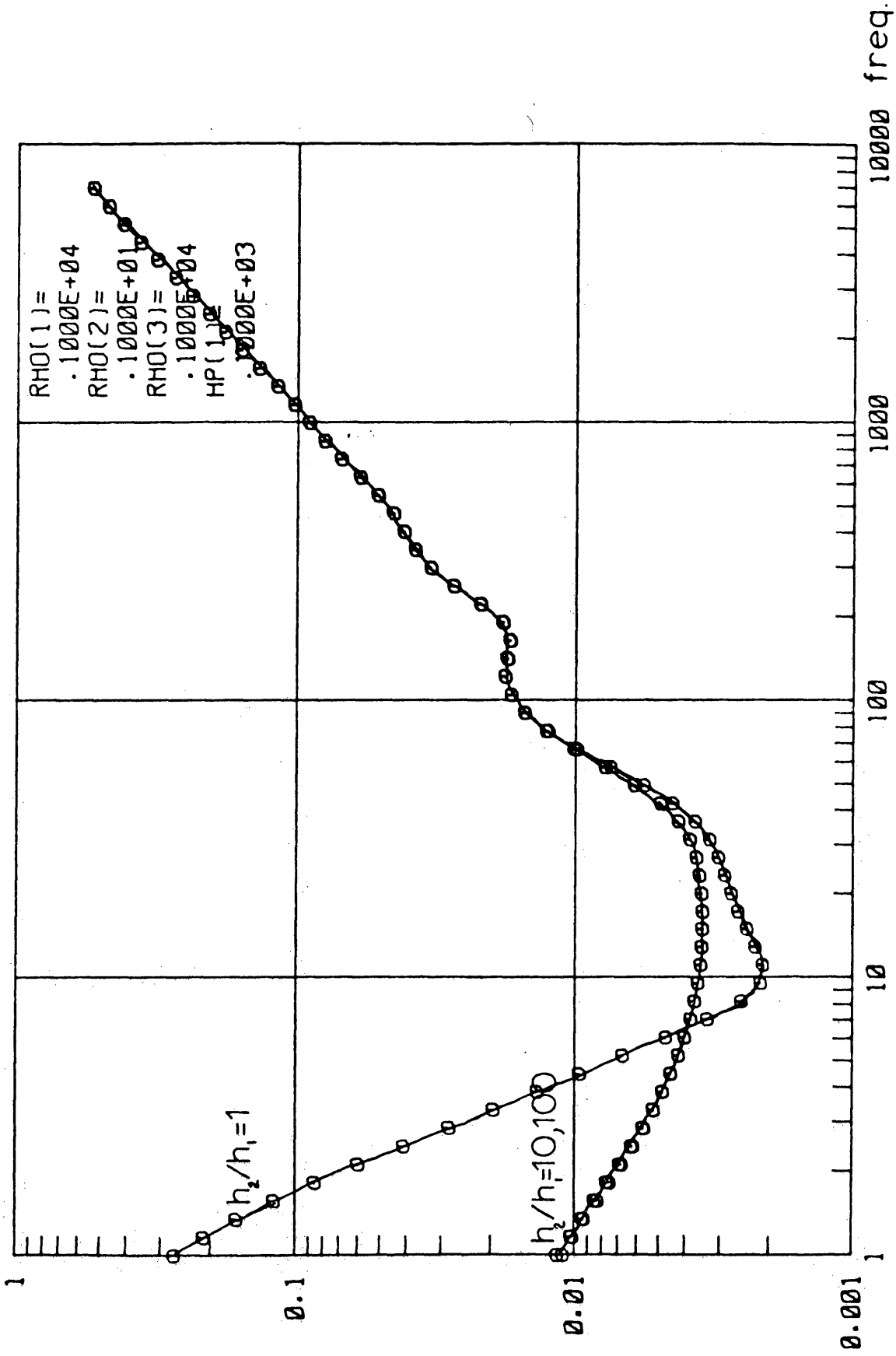


IH/H1I 3 LAYERS R = .500E+05

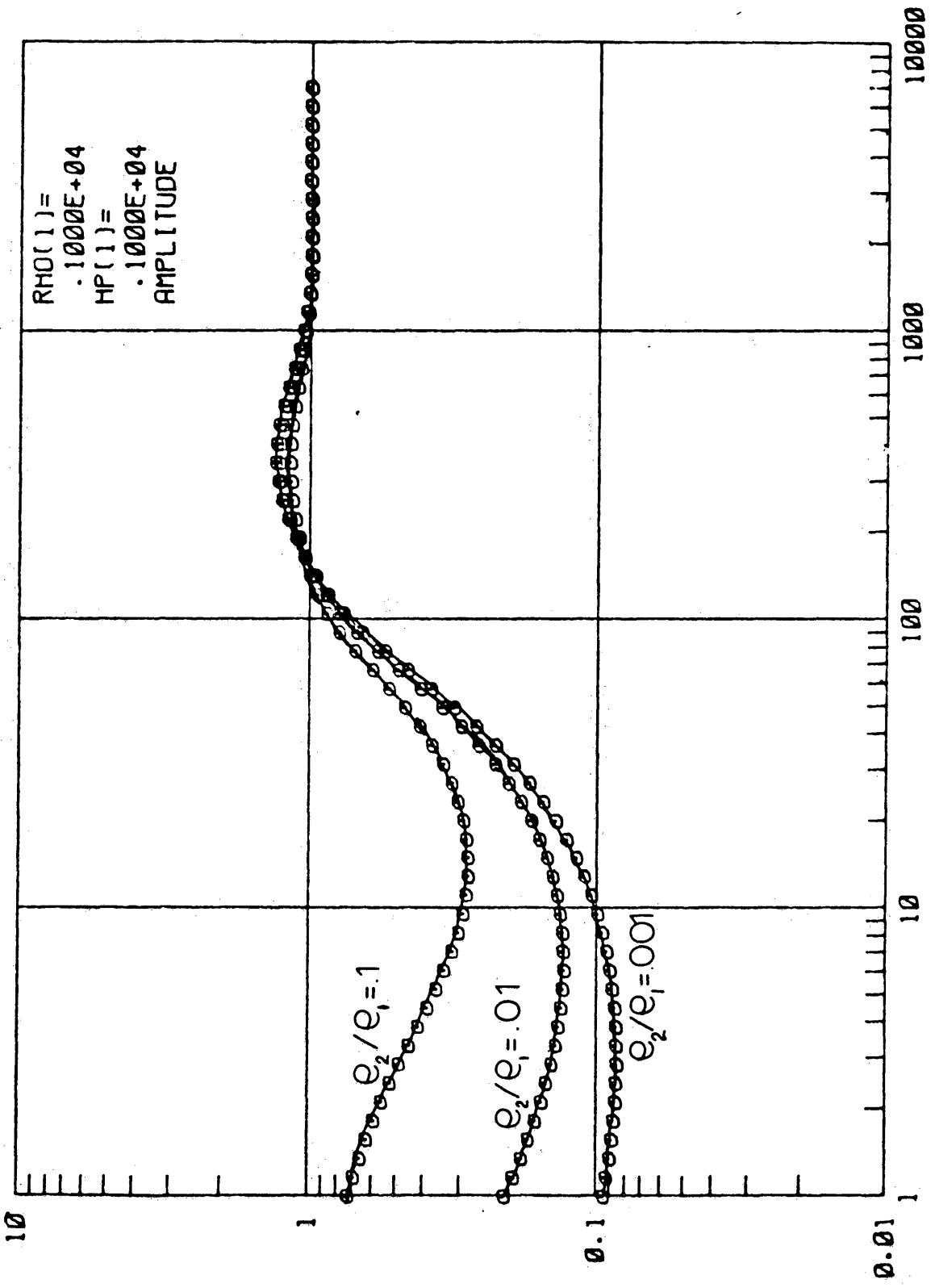
RHO(1)=  
 .1000E+04  
 RHO(2)=  
 .1000E+01  
 RHO(3)=  
 .1000E+04  
 HP(1)=  
 .1000E+03



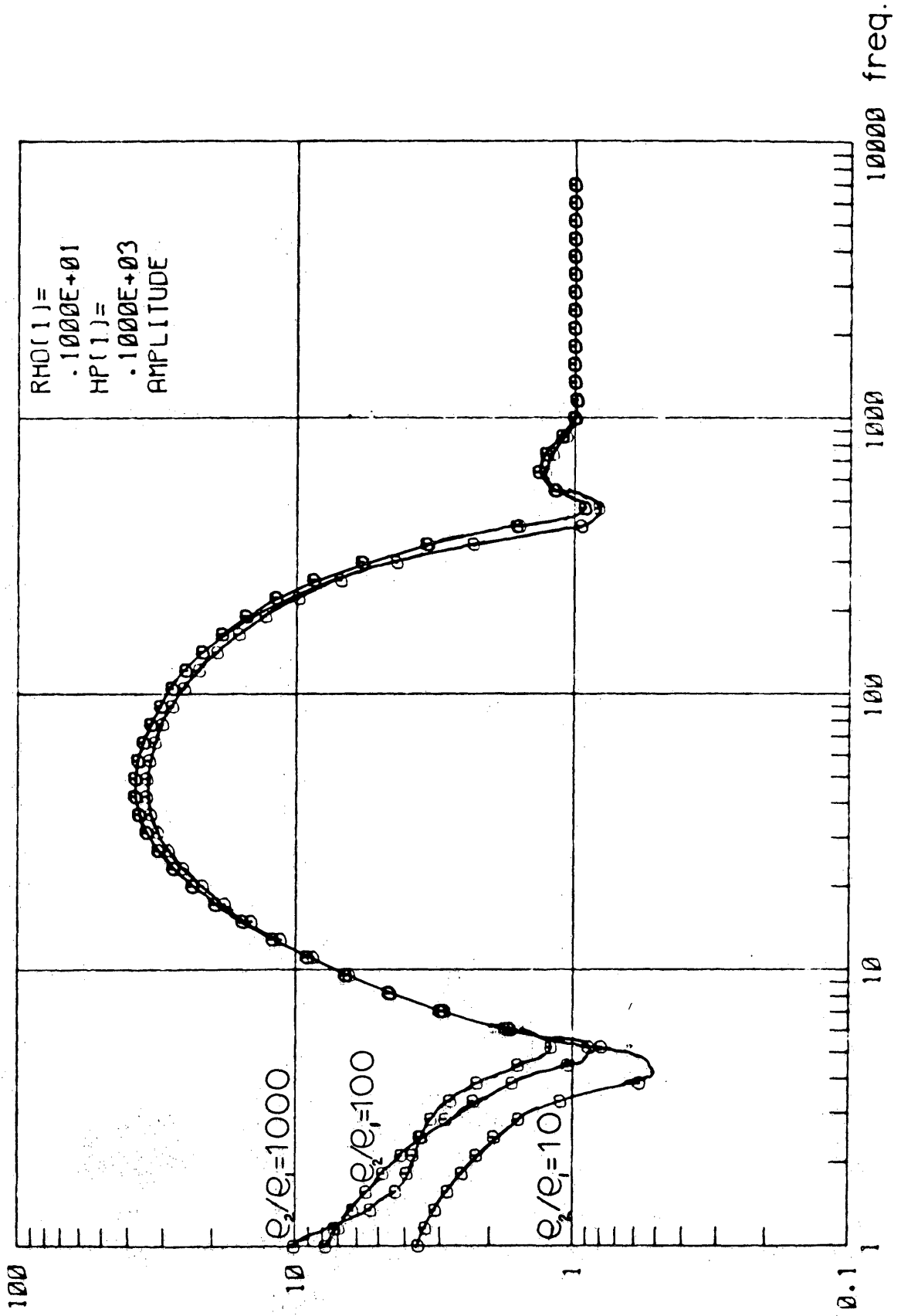
|H/H1| 3 LAYERS R= .500E+05



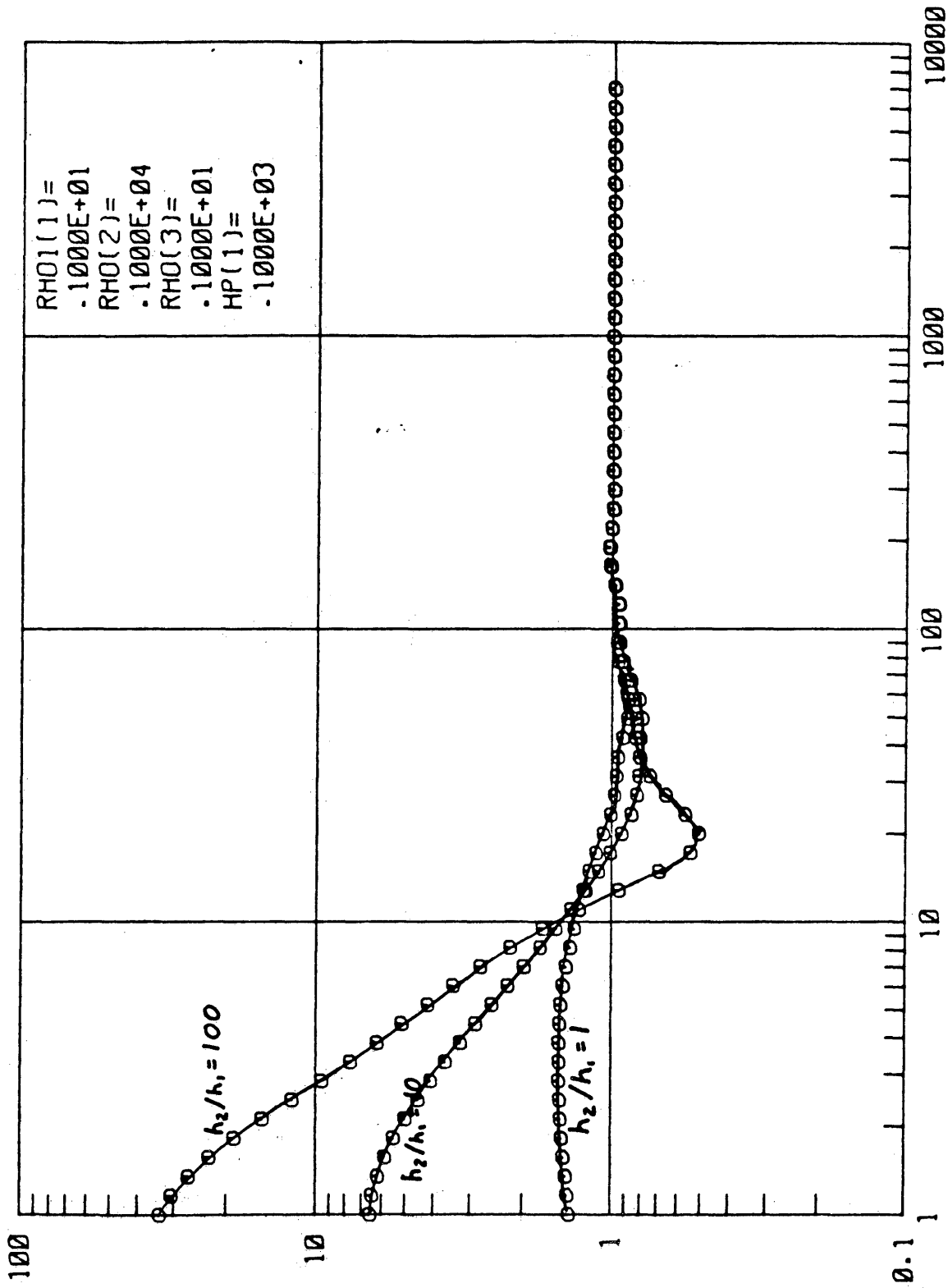
1H/H11 3 LAYERS R = .100E+05



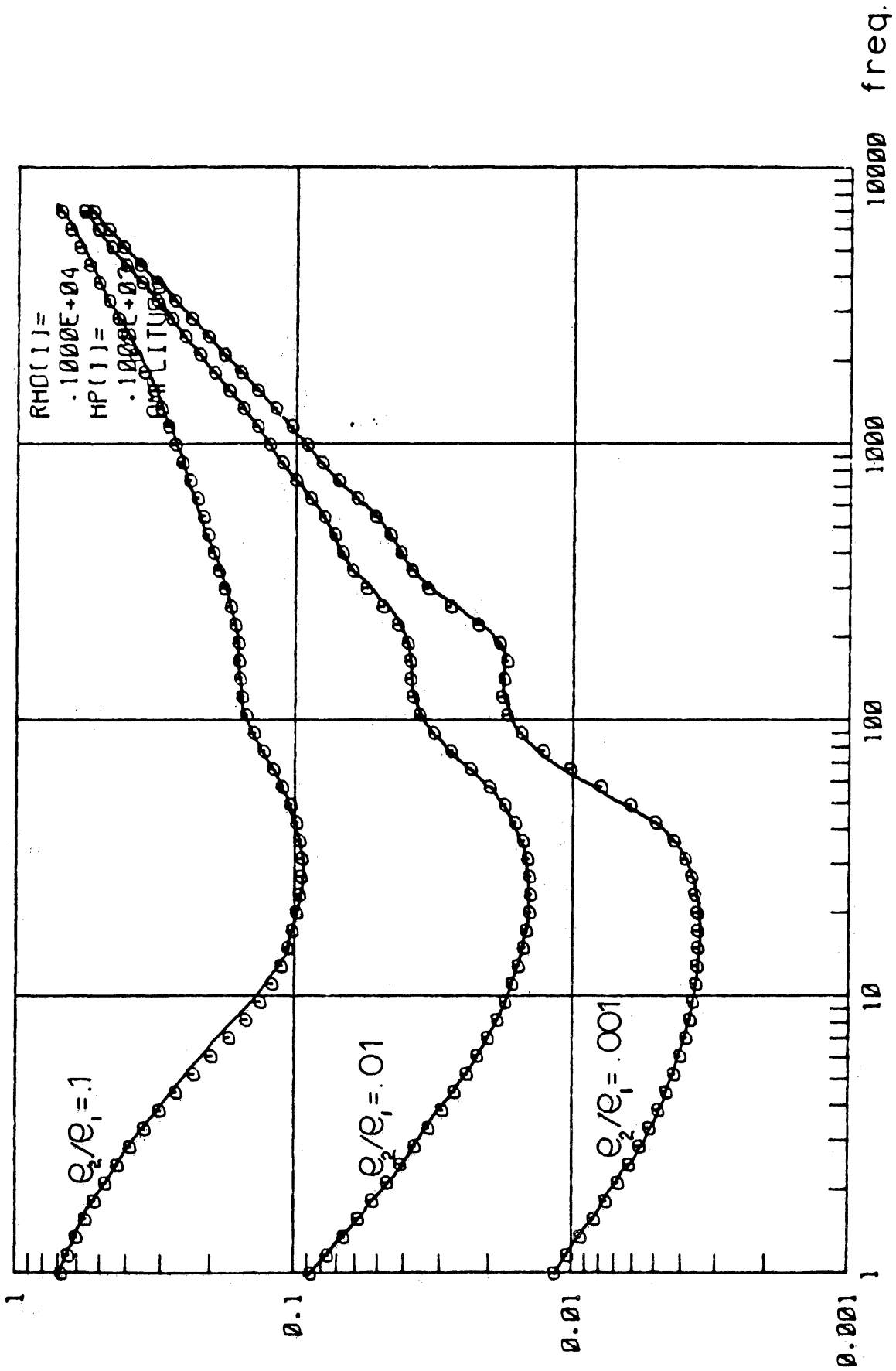
H/H1 2 LAYERS R = .100E+05



H/H11 2 LAYERS R = .500E -



IH/H11. 3 LAYERS R= .100E+05



IH/H11 2 LAYERS R = .100E+05

APPENDIX B  
COMPUTER PROGRAMS

```

C***** PROGRAM EMFWD *****
C
C   THIS PROGRAM CALCULATES THE VERTICAL H FIELD, THE X-COMPONENT
C   OF THE ELECTRIC FIELD, THE Y-COMPONENT OF THE ELECTRIC FIELD,
C   AND THE COMBINED HORIZONTAL ELECTRIC FIELD COMPONENTS (EC=
C   SIN(2*AVG)EX-COS(2*AVG)EY) FOR A 1,2,OR 3 LAYERED EARTH.  THE SOURCE
C   IS A HORIZONTAL CURRENT DIPOLE
C
C   FIELDS FOR A HOMOGENEOUS HALF-SPACE (1 LAYER) MAY BE COMPUTED
C   BY USING A TWO LAYER MODEL WITH RH0(2)=RH0(1).  MODELS
C   FOR MORE THAN THREE LAYERS MAY BE COMPUTED SIMPLY BY INCREASING
C   THE DIMENSIONS OF THE ARRAYS.
C
C   PROGRAM EMFWD WAS DEVELOPED ON A DEC PDP-10 COMPUTER AT THE
C   COLORADO SCHOOL OF MINES.  FORTRAN PROGRAMING LANGUAGE IS USED
C   THROUGHOUT
C
C
C           JEFF DANIELS
C   COLORADO SCHOOL OF MINES
C           FEBRUARY 1974
C
C***** VARIABLES *****
C
C NFW= THE NUMBER OF FREQUENCIES TO BE CALCULATED
C R= SOURCE-RECEIVER SEPERATION
C AVG=ANGLE, IN DEGREES, DEFINING SOURCE-RECEIVER ORIENTATION
C FF=FREQUENCY IN HERTZ
C F=ANGULAR FREQUENCY
C N= THE NUMBER OF LAYERS
C HP(I)= LAYER THICKNESS
C RH(I)= LAYER RESISTIVITY
C CH(I)= COEFFICIENTS FOR CALCULATING THE J1 HANKEL TRANSFORM
C YH(I)= ABSCISSA VALUES FOR CH(I)
C CE(I)= COEFFICIENTS FOR CALCULATING THE J0 HANKEL TRANSFORM
C YE(I)= ABSCISSA VALUES FOR CALCULATING CE(I)
C CI = SOURCE-DIPOLE CURRENT
C DS= SOURCE-DIPOLE LENGTH
C D = NORMALIZED THICKNESSES
C RK= NORMALIZED RESISTIVITIES
C DEL= MODIFIED WAVE NUMBER
C IIN= INPUT AREA
C IOUT= OUTPUT AREA
C
C
C*****
C   COMPLEX H, E, EX, EY
C   DIMENSION HPP(3),HP(3)

```

COMMON /QC/D(3),RK(3),RH(3),V,DEL,R

COMMON /CV/CI,DS,ANG,F,X,YH(48),CH(48),YE(61),CE(61),TM

DATA

&YH/- 4.5307316E 0,- 4.3004731E 0,- 4.0702146E 0,- 3.8399561E 0,  
 & -3.6096976E 0,- 3.3794391E 0,- 3.1491806E 0,- 2.9189221E 0,  
 & -2.6886636E 0,- 2.4584051E 0,- 2.2281466E 0,- 1.9978881E 0,  
 & -1.7676296E 0,- 1.5373711E 0,- 1.3071126E 0,- 1.0768541E 0,  
 & -8.4659563E-1,- 6.1633713E-1,- 3.8607863E-1,- 1.5582013E-1,  
 & 7.4438369E-2, 3.0469687E-1, 5.3495537E-1, 7.6521387E-1,  
 & 9.9547237E-1, 1.2257309E 0, 1.4559894E 0, 1.6862479E 0,  
 & 1.9165064E 0, 2.1467649E 0, 2.3770234E 0, 2.6072819E 0,  
 & 2.8375404E 0, 3.0677989E 0, 3.2980574E 0, 3.5283159E 0,  
 & 3.7585744E 0, 3.9888329E 0, 4.2190914E 0, 4.4493499E 0,  
 & 4.6796084E 0, 4.9098669E 0, 5.1401254E 0, 5.3703839E 0,  
 & 5.6006424E 0, 5.8309009E 0, 6.0611594E 0, 6.2914179E 0/

DATA

&CH/3.1010561E-6, 1.8802098E-5, 5.4819540E-5, 9.2891602E-6,  
 & 1.5523239E-4, 3.0344652E-5, 3.5338744E-4, 1.4793002E-4,  
 & 7.7342377E-4, 5.3570357E-4, 1.7170605E-3, 1.6387239E-3,  
 & 3.9247683E-3, 4.5796508E-3, 9.2111468E-3, 1.2130467E-2,  
 & 2.1938415E-2, 3.0853660E-2, 5.1973594E-2, 7.4661566E-2,  
 & 1.1775455E-1, 1.6353574E-1, 2.3127545E-1, 2.7368461E-1,  
 & 2.8059285E-1, 1.2875840E-1,- 1.5380437E-1,- 4.5659951E-1,  
 & -3.6077766E-2, 4.2985683E-1,- 2.1506075E-1,- 2.3624312E-2,  
 & 8.9316746E-2,- 7.4344203E-2, 4.8572965E-2,- 3.0083872E-2,  
 & 1.8846544E-2,- 1.2158687E-2, 8.0708759E-3,- 5.4706275E-3,  
 & 3.7554604E-3,- 2.5929707E-3, 1.7909426E-3,- 1.2320277E-3,  
 & 8.4095286E-4,- 5.6749747E-4, 3.7718405E-4,- 1.5891835E-4/

DATA

&YE/- 6.8348046E 0,- 6.6045461E 0,- 6.3742876E 0,- 6.1440291E 0,  
 & -5.9137706E 0,- 5.6835121E 0,- 5.4532536E 0,- 5.2229951E 0,  
 & -4.9927366E 0,- 4.7624781E 0,- 4.5322196E 0,- 4.3019611E 0,  
 & -4.0717026E 0,- 3.8414441E 0,- 3.6111856E 0,- 3.3809271E 0,  
 & -3.1506686E 0,- 2.9204101E 0,- 2.6901516E 0,- 2.4598931E 0,  
 & -2.2296346E 0,- 1.9993761E 0,- 1.7691176E 0,- 1.5338591E 0,  
 & -1.3086006E 0,- 1.0783421E 0,- 8.4808353E-1,- 6.1782503E-1,  
 & -3.8756658E-1,- 1.5730308E-1, 7.2950416E-2, 3.0320892E-1,  
 & 5.3346742E-1, 7.6372592E-1, 9.9398442E-1, 1.2242429E 0,  
 & 1.4545014E 0, 1.6847599E 0, 1.9150184E 0, 2.1452769E 0,  
 & 2.3755354E 0, 2.6057939E 0, 2.8360524E 0, 3.0663109E 0,  
 & 3.2965694E 0, 3.5268279E 0, 3.7570864E 0, 3.9873449E 0,  
 & 4.2176034E 0, 4.4478619E 0, 4.6781204E 0, 4.9083789E 0,  
 & 5.1386374E 0, 5.3688959E 0, 5.5991544E 0, 5.8294129E 0,  
 & 6.0596714E 0, 6.2899299E 0, 6.5201884E 0, 6.7504469E 0,  
 & 6.9807054E 0/

DATA

&CE/7.3260937E-4, 5.6326423E-4, 1.3727237E-4, 7.5331222E-4,  
 & 3.5918326E-4, 1.0500608E-3, 7.1530982E-4, 1.5160070E-3,  
 & 1.2841617E-3, 2.2497985E-3, 2.1906186E-3, 3.4076782E-3,  
 & 3.6321245E-3, 5.2376028E-3, 5.9212519E-3, 8.1315877E-3,  
 & 9.5527062E-3, 1.2708615E-2, 1.5305589E-2, 1.9941086E-2,

```

& 2. 439 662 6E-2, 3. 1333652E-2, 3. 868 3065E-2, 4. 912799 3E-2,
& 6. 0824806E-2, 7. 631 4344E-2, 9. 3928 346E-2, 1. 1545027E-1,
& 1. 38 68 663E-1, 1. 62438 47E-1, 1. 811 4332E-1, 1. 8424433E-1,
& 1. 5556741E-1, 6. 8592481E-2, -8. 8339029E-2, -2. 8819226E-1,
& -3. 5565260E-1, -5. 6288 677E-2, 4. 818 69 42E-1, -5. 1516453E-2,
& -2. 6102989E-1, 2. 141 6490E-1, -9. 449068 7E-2, 2. 6196370E-2,
& -5. 1097328E-4, -6. 60329 48E-3, 7. 5193619E-3, -6. 78 54344E-3,
& 5. 8044372E-3, -4. 935 439 4E-3, 4. 2323106E-3, -3. 6733648E-3,
& 3. 2266260E-3, -2. 8649 137E-3, 2. 5677680E-3, -2. 3202655E-3,
& 2. 1115187E-3, -1. 9334662E-3, 1. 7800248E-3, -1. 6465436E-3,
& 1. 3468317E-3/

```

```

14  FORMAT(1X, I2, 1X, 'LAYERS', 2X, 'SOURCE CURRENT= ',
&F8. 3, 2X, 'SOURCE LENGTH= ', F8. 2, /, 'ANGLE= ', F8. 2, 2X,
&'SOURCE-RECEIVER SEPERATION= ', E12. 4, /)
15  FORMAT(3X, 'LAYER', 3X, 'RESISTIVITY', 3X, 'THICKNESS')
16  FORMAT(1X, I5, 4X, E11. 5, 3X, E10. 4)
10  FORMAT(/, 3X, 'FREQUENCY', 4X, 'H(REAL)', 6X, 'H(IMAG)', 6X, 'EC
&(REAL)', 5X, 'EC(IMAG)', /, 3X, 'FREQUENCY', 4X, 'EX(REAL)',
&5X, 'EX(IMAG)', 5X, 'EY(REAL)', 5X, 'EY(IMAG)', //)
12  FORMAT(5(1X, E12. 6))
5   FORMAT(1F)
1   FORMAT(2F, 2I)
28  FORMAT(2F)
3   FORMAT(2F)
    IIN=3
    IOUT=8
    READ(IIN, 1) AVG, R, N, VFW
    DØ 11 LT=1, N
11  READ(IIN, 28) HP(LT), RH(LT)
    READ(IIN, 3) CI, DS
    WRITE(IOUT, 14) N, CI, DS, AVG, R
    WRITE(IOUT, 15)
    WRITE(IOUT, 16) (LT, RH(LT), HP(LT), LT=1, N)
    VB=2*N
    TM=12. 566371E-7
    AVG=(3. 1415927/180.)*AVG
    DØ 2 JJ=1, N
2   RK(JJ)=RH(1)/RH(JJ)
    WRITE(IOUT, 10)
    DØ 20 I=1, VFW
    READ(IIN, 5) FF
    F=6. 2831853*FF
    DEL=SQRT(2.*RH(1)/(TM*F))
    DØ 6 JJ=1, N
6   D(JJ)=2.*HP(JJ)/DEL
    X=ALOG(R/DEL)
    CALL FV(E, H, EX, EY)
    WRITE(IOUT, 12) FF, H, E
    WRITE(IOUT, 12) FF, EX, EY
20  CONTINUE
    END

```

## SUBROUTINE FV(EC, H, EX, EY)

```

C
C*****SUBROUTINE FVALUE *****
C THIS SUBROUTINE CALCULATES THE CONVOLUTION SUMS FOR EX,
C EY, EC, AND HZ ELECTROMAGNETIC FIELD COMPONENTS
C
C***** VARIABLES *****
C
C SUM= CONVOLUTION SUM FOR HZ
C SUEY1, SUEY0, SUEX1, SUEX0 AS DEFINED ON PAGE 10 OF THESIS
C SUEY1= CONVOLUTION SUM FOR EY1
C SUEY0= CONVOLUTION SUM FOR EY0
C SUEX1= CONVOLUTION SUM FOR EX1
C SUEX0= CONVOLUTION SUM FOR EX0
C H1= CLOSED-FORM HOMOGENEOUS HALF-SPACE EXPRESSION FOR HZ
C EX= X-COMPONENT OF THE ELECTRIC FIELD FOR A LAYERED EARTH
C EY= Y-COMPONENT OF THE ELECTRIC FIELD FOR A LAYERED EARTH
C EX1= CLOSED-FORM EXPRESSION FOR A HOMOGENEOUS HALF-SPACE
C      (X-COMPONENT)
C EY1= CLOSED-FORM EXPRESSION FOR A HOMOGENEOUS HALF-SPACE
C      (Y-COMPONENT)
C E= CLOSED-FORM EXPRESSION FOR A HOMOGENEOUS HALF-SPACE
C   ( COMBINED ELECTRIC FIELD)
C EC= COMBINED ELECTRIC FIELD FOR A LAYERED EARTH
C F1= LAYERED EARTH CORRECTION FACTOR
C Z1= LAYERED EARTH CORRECTION FACTOR
C V1= NORMALIZED FIRST LAYER PSEUDO WAVE NUMBER
C
C
C*****
      COMMON /CV/ CI, DS, AVG, F, X, YH(48), CH(48), YE(61), CE(61), TM
      COMMON /QC/ D(3), R(3), RH(3), V, DEL, R
      COMPLEX H, V1, F1, SUM, H1, Z1, E1, EC, C1, C2,
&EX1, EX0, EY1, EY0, SUEX1, SUEY1, SUEX0, SUEY0, EX, EY
      SUEX1=CMPLX(0.0, 0.0)
      SUEY1=CMPLX(0.0, 0.0)
      SUEX0=CMPLX(0.0, 0.0)
      SUEY0=CMPLX(0.0, 0.0)
      SUM=CMPLX(0.0, 0.0)
      BB=EXP(X)
      C1=CMPLX(0.0, 1.0)
      C2=CMPLX(1.0, 1.0)
C*****
CALCULATE THE J1 HANKEL TRANSFORMS
C*****
      DO 8 J=1, 48
      Y=EXP(-(X-YH(J)))
      CALL CALC(F1, Z1, V1, Y)
      EX1=C1*V1*(Z1-1)+2.*V1*(1-F1)/((Y+V1*F1)*(Y+V1))
      SUEX1=SUEX1+CH(J)*EX1
      H=Y*Y*V1*(1.-F1)

```

```

H=H/((Y+V1*F1)*(Y+V1))
SUM=SUM+CH(J)*H
8 CONTINUE
SUEY1=SUEX1
H=(1/BB)*CI*DS*SIN(ANG)*SUM/(6.2831853*DEL*DEL)
R2=R*R
H1=3.-(3.+3.0*C2*BB+2.*C1*BB**2)*CEXP(-C2*BB)
H1=(-C1*CI*DS*DEL*DEL*SIN(ANG)*H1)/(12.566371*R2*R2)
12 FORMAT(2(1X,E14.8))
H=H+H1
SUM=CMPLX(0.0,0.0)
C*****
CALCULATE THE JO HAVKEL TRANSFORMS
C*****
DØ 4 I=1,61
Y=EXP(-(X-YE(I)))
CALL CALC(F1,Z1,V1,Y)
EY0=Y*(C1*V1*(Z1-1.)+2.*V1*(1.-F1)/((Y+V1)*(Y+V1*F1)))
EX0=(CØS(ANG)**2/DEL)*EY0-(1/DEL)*2.*Y*V1*(1.-F1)/
&((Y+V1*F1)*(Y+V1))
SUEX0=SUEX0+CE(I)*EX0
SUEY0=SUEY0+CE(I)*EY0
EC=C1*V1*Y*(Z1-1.0)/2.0
EC=EC-Y*V1*(1.0-F1)/((Y+V1*F1)*(Y+V1))
4 SUM=SUM+CE(I)*EC
EC=C1*TM*F*CI*DS*SIN(2.*ANG)*SUM/(12.566371*DEL*BB)
E1=-.500*(1+C2*BB)*CEXP(-C2*BB)
E1=CI*DS*RH(1)*SIN(2.*ANG)*E1/(6.2831853*R*R2)
EC=E1+EC
EX=(C1*TM*F*CI*DS/(12.566371*BB))*(SUEX0-CØS(2.*ANG)
&*SUEX1/R)
EY=(C1*TM*F*CI*DS/(12.566371*BB))*SIN(2.*ANG)*(-SUEY1/R
&+SUEY0/(2.*DEL))
EX1=CI*DS*RH(1)*((3*CØS(ANG)**2-2)+(1.+(1+C1)*BB)*CEXP(-C2*BB)
&)/(6.2831853*R*R2)
EY1=3.*RH(1)*CI*DS*SIN(ANG)*CØS(ANG)/(6.2831853*R*R2)
EX=EX+EX1
EY=EY+EY1
RETURN
END

```

```

SUBROUTINE CALC(F1,Z1,V1,Y)
C*****
C SUBROUTINE CALC CALCULATES THE LAYERED EARTH CORRECTION FACTORS
C (F1 AND Z1) FOR UP TO 3 LAYERS*****
COMMON /QC/ D(3),RK(3),RH(3),N,DEL,R
COMPLEX V1,F1,V2,AEX,Z1
Y2=Y*Y
T=2.*RK(N)
V2=CSQRT(CMPLX(Y2,T))
DO 9 LL=2,N
IF(LL.GT.2) GO TO 4
F1=CMPLX(1.0,0.0)
Z1=CMPLX(1.0,0.0)
4 I=N-LL+1
DD=D(I)
T=2.*RK(I)
V1=CSQRT(CMPLX(Y2,T))
AEX=(1.-CEXP(-V1*DD))/(1.+CEXP(-V1*DD))
F1=(V2*F1+V1*AEX)/(V1+V2*F1*AEX)
Z1=(V2*RH(I+1)*Z1+V1*RH(I)*AEX)/(V1*RH(I)+V2*RH(I+1)*Z1*AEX)
9 V2=V1
RETURN
END

```



```

COMMON /CV/ CI, DS, AVG, FX, XX, TM
COMMON /QC/ DD(3), RKK(3), RHH(3), V, DEL, RR
DIMENSION X(25), Y(25), B(7), WTS(25),
&BDB(7), ATA(7, 7), IFXE(6), HP(3), RH(3)
REAL LAM
IV=1

```

```

C*****
C READ IN THE DATA UNDER THE LOGICAL UNIT "IV"
C*****
  READ(IV, 1) AVG, R, CI, DS, IE, V, VCOMP
 7  FORMAT(E)
 1  FORMAT(4F, 3I)
  READ(IV, 9) (X(I), Y(I), WTS(I), I=1, VFW)
 9  FORMAT(3E12.6)
 2  FORMAT(2F)
  AVG=(3.1415927/180.)*AVG
  VB=2*V
  READ(IV, 8) (IFXE(J), J=1, VB)
 8  FORMAT(I6)
  READ(IV, 2) (HP(LT), RH(LT), LT=1, V)
  READ(IV, 7) YVORM
  DO 4 JJ=1, VFW
 4  Y(JJ)=Y(JJ)/YVORM
  VPAR=VB
  B(1)=R
  DO 3 JCT=1, V
  VCT=2*JCT
  B(VCT)=RH(JCT)
  B(VCT+1)=HP(JCT)
 3  CONTINUE
  CALL NVNLS2(VPAR, IFXE, IE, X, Y, WTS, B, VCOMP, YVORM)
  END

```

```

SUBROUTINE V0VLS2 (NPAR,IFXE,IE,X,Y,WTS,B,VCOMP,YNORM)
C***** SUBROUTINE V0VLS2 *****
C THIS SUBROUTINE CONTROLS THE LEAST-SQUARES CURVE FITTING
C ALGORITHM AS OUTLINED BY MARQUARDT (1963)
C
C
C THIS IS A MODIFIED VERSION OF THE ORIGINAL
C GEOPHYSICS LEAST-SQUARES PROGRAM WRITTEN BY JORGE
C PARRA.
C
C
C*****LEAST SQUARES SUBROUTINE*****
C NVAR=# OF UNFIXED PARAMETERS
C NPAR=# OF UNKNOWN PARAMETERS OF F > LINOUT
C NPTS=# OF OBSERVATION POINTS
C IW=OUTPUT (FLAG:-1=NO OUTPUT,0,1,2=OUTPUT)
C FAIL=OUTPUT(0,1,2=CONVERGENCE;3,4,5=NO CONVERGENCE)
C MAXITR=THE MAXIMUM # OF ITERATIONS ALLOWED
C IFXE=ACCURACY(IFXE(K)=1,THE KTH PARAMETER IS ALLOWED TO VARY;
C IFXE(K)=0, THE KTH PARAMETER IS FIXED AT ITS INPUT VALUE)
C NEXI=0
C LINMIN=0
C X(IE,2)=DOUBLY DIMENSIONED ARRAY HOLDING THE OBSERVATIONS
C (INDEPENDENT VARIABLES) AS ROWS OF X
C Y(IE)= ARRAY HOLDING THE DEPENDENT VARIABLE OBSERVATIONS
C WTS=ARRAY HOLDING THE RELIABILITY WEIGHTS FOR EACH POINT
C /IF WTS=0 THE WEIGHT =1.0/
C ATA=A* MATRIXFROM MARQUARDT'S FORMULATION
C B=HOLDS THE INITIAL ESTIMATES OF THE PARAMETERS. IF IFAIL=0
C OR 1, B HOLDS THE CONVERGED PARAMETER SET
C BDB= THE FINAL CORRECTION VECTOR TO THE CONVERGED PARAMETERS
C IP=10**-5
C LAM=MARQUARDT'S DAMPING CONSTANT
C PHI=(Y(OBS)-Y(MODEL))**2
C DAC=.05 / PARABOLIC MINIMIZATION
C FVALUE=SUBROUTINE THAT CALCULATES THE MODEL
C PARVAL=SUBROUTINE THAT CALCULATES THE FIRST DERIVATIVE OF THE
C MODEL
C ANS=CORRECTION VECTOR AT THE POINTS OF LINEAR MINIMIZATION
C *****
DIMENSION ATF(7),ANS(7),ANS1(7),EX(7)
DIMENSION X(25),Y(25),B(7),BDB(7),WTS(25),ATA(7,7)
REAL LAM,LD0,L10,NORM
DIMENSION INX(7),IFXE(7)
INTEGER FAIL
1003 FORMAT(4(2X,E14.8))
C*****
C SET THE CONSTANT VALUES
C*****
LAM= 0.01
NPTS=IE

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```

ID=VPAR
MAXITR=6
EP=1.E+0
VEXT=0
GAMA0=45.
  IW=3
ITER=0
TOLR=1.E-8
IOUT=2
IR=3
C   DETERMINE THE PARAMETERS TO BE VARIED.
C
VVAR=0
DØ 10 I=1,VPAR
BDR(I)=B(I)
AVS(I)=0.E0
AVSI(I)=0.E0
IF(IFXE(I).LE.0) GØ TØ 10
VVAR=VVAR+1
IX(VVAR)=I
10  CONTINUE
3001 FORMAT(I)
IF(IW.GE.0) WRITE(IOUT,11700) VPAR,VVAR,VPTS,MAXITR,IW,LINMIN,EP,
*DAC
C
C   EXAMINE THE WTS ARRAY
C
VØRM=0.E0
DØ 40 I=1,VPTS
IF(WTS(I)) 20,30,40
20  WTS(I)=1.E0/WTS(I)**2
GØ TØ 40
30  WTS(I)=1.E0
40  VØRM=VØRM+WTS(I)
    VØRM=FLOAT(VPTS)/VØRM
IF(VØRM.EQ.1.E0) GØ TØ 60
DØ 50 I=1,VPTS
50  WTS(I)=WTS(I)*VØRM
C
C   CALCULATE THE INITIAL SUM OF SQUARES
C
60  PHI=0.E0
IF(IW.GT.0) WRITE(IOUT,10200)
DØ 70 I=1,VPTS
VPVFC=1
CALL EMCALC (X,B,IE,I,EX,FV,IFXE,VPVFC,VCØMP,YVØRM)
SS=Y(I)-FV
IF(IW.GT.0) WRITE(IOUT,10000) X(I),Y(I),FV,SS
70  PHI=PHI+SS**2*WTS(I)
C*****
C   START THE DAMPED GAUSSIAN PROCEDURE

```

```

C*****
C
80  ITER=ITER+1
   IF(IW.GT.0) WRITE(IOUT,11800)
   IF(ITER.GT.MAXITER) GO TO 380
   DO 90 I=1,NPAR
90  BDB(I)=B(I)
   DO 100 I=1,NVAR
   ATF(I)=0.E0
   DO 100 J=I,NVAR
100  ATA(I,J)=0.E0
C*****
C      GENERATE THE ATA AND ATF ARRAYS
C*****
   DO 110 I=1,NPTS
   NPVFC=1
   CALL EMCALC(X,B,IE,I,EX,FV,IFXE,NPVFC,NCOMP,YNORM)
   NORM=WTS(I)*(Y(I)-FV)
   NPVFC=0
   CALL EMCALC(X,B,IE,I,EX,FV,IFXE,NPVFC,NCOMP,YNORM)
   DO 110 L=1,NVAR
   J=INX(L)
   ATF(L)=ATF(L)+NORM*EX(J)
   DO 110 M=L,NVAR
   K=INX(M)
   ATA(L,M)=ATA(L,M)+EX(J)*EX(K)*WTS(I)
110  CONTINUE
   IF(IW.GT.0) WRITE(IOUT,11500) ITER,PHI,(B(I),I=1,NPAR)
C*****
C      PERFORM A LOCAL SCALING ON THE ATA MATRIX TO AID CALCULATIONS.
C*****
   DO 120 I=1,NVAR
   IF(ATA(I,I).EQ.0.D0) GO TO 410
120  EX(I)=SQRT(ATA(I,I))
   DO 140 I=1,NVAR
   ATF(I)=ATF(I)/EX(I)
   DO 140 J=I,NVAR
   IF(I.EQ.J) GO TO 130
   ATA(I,J)=ATA(I,J)/(EX(I)*EX(J))
   GO TO 140
130  ATA(I,I)=1.E0
140  CONTINUE
C*****
C      DETERMINE A VALID LAMDA FOR THE SCALED PARTIAL MATRIX.
C      GAMA=DIRECTION VECTOR ANGLE
C      PHI1=SS ERROR FOR NEW LAMBDA
C*****
   FAC=1.E+0
   CALL NEWLAM(ATA,LAM,BDB,ATF,ANS,EX,GAMA,FVALUE,PHI1,X,Y,WTS,B,
 *ID,IE,NVAR,NPTS,NPAR,IFAL,INX,IFXE,IR,NCOMP,YNORM)
   IF(IW.GT.0) WRITE(IOUT,10600) LAM,PHI1,GAMA

```

```

C*****
C EP= RELATED TO THE FIRST GUESS (CONTROLS LAMBDA SO THAT
C LAMBDA DOESN'T GET TOO LARGE) EXAMPLE: EP=.5 THEN THE
C GREATEST CHANGE THAT THE PARAMETERS CAN HAVE IS 50% OF
C THE ORIGINAL FIRST GUESS
C*****
      IF(IFAL-1) 150, 390, 400
150   DO 160 I=1, NVAR
      IF(ABS(ANS(I)).GT.(TOLR+EP*ABS(B(INX(I))))) GO TO 180
160   CONTINUE
      IF(PHI.LT.PHI1) GO TO 440
      DO 170 I=1, NVAR
170   B(INX(I))=B(INX(I))+ANS(I)
      PHI =PHI1
      GO TO 440
180   IF(PHI1.GE.PHI) GO TO 220
      IF(LAM.LE.TOLR) GO TO 320
      DO 1213 JPI=1, NVAR
1213  B(JPI)=BDB(JPI)
      PHI=PHI1
      LD0=LAM/10.0E0
      CALL NEWLAM(ATA,LD0,BDB,ATF,ANS1,EX,GAMA,FVALUE,PHI2,X,Y,WTS,B,
      *ID,IE,NVAR,NPTS,VPAR,IFAL,INX,IFXE,IR,VCOMP,YNORM)
      IF(IW.GT.0) WRITE(IOUT,10700) LD0,PHI2,GAMA
      IF(IFAL-1) 190, 390, 400
190   IF(PHI2.GE.PHI1) GO TO 320
200   LAM=LD0
      DO 1207 JPI=1, NVAR
1207  B(JPI)=BDB(JPI)
      PHI=PHI2
      DO 210 I=1, NVAR
210   ANS(I)=ANS1(I)
      PHI1=PHI2
      GO TO 320
220   LD0=LAM/10.0E0
      CALL NEWLAM(ATA,LD0,BDB,ATF,ANS1,EX,GAMA,FVALUE,PHI2,X,Y,WTS,B,
      *ID,IE,NVAR,NPTS,VPAR,IFAL,INX,IFXE,IR,VCOMP,YNORM)
      IF(IW.GT.0) WRITE(IOUT,10700) LD0,PHI2,GAMA
      IF(IFAL-1) 230, 390, 400
230   IF(PHI2.LT.PHI) GO TO 200
      L10=LAM
240   L10=L10*10.E0
      CALL NEWLAM(ATA,L10,BDB,ATF,ANS,EX,GAMA,FVALUE,PHI3,X,Y,WTS,B,
      *ID,IE,NVAR,NPTS,VPAR,IFAL,INX,IFXE,IR,VCOMP,YNORM)
      IF(L10.GT.1.E+6) STOP
      IF(IW.GT.0) WRITE(IOUT,10900) L10,PHI3,GAMA
      IF(IFAL-1) 250, 390, 400
250   IF(PHI3.GE.PHI.AND.L10.GT.1.E+3) GO TO 360
      IF(PHI3.GE.PHI) GO TO 260
      PHI1=PHI3
      PHI=PHI3

```

```

      DØ 1208 JPI=1, VVAR
1208  B(JPI)=BDB(JPI)
      LAM=L10
      GØ TØ 320
260  IF(GAMA.GE.GAMA0) GØ TØ 240
      FAC=FAC/2.E0
      DØ 270 I=1, VVAR
270  ANS(I)=ANS(I)/2.E0
      DØ 280 I=1, VVAR
      IF(ABS(ANS(I)).GT.(TØLR+EP*ABS(B(INX(I)))))) GØ TØ 290
280  CONTINUE
      GØ TØ 430
290  DØ 300 I=1, VVAR
300  BDB(INX(I))=B(INX(I))+ANS(I)
      PHI3=0.D0
      DØ 310 I=1, VPTS
      VPVFC=1
      CALL EMCALC(X, BDB, IE, I, EX, FV, IFXE, VPVFC, VCOMP, YVØRM)
310  PHI3=PHI3+(Y(I)-FV)**2*WTS(I)
      IF(IW.GT.0) WRITE(IØUT,10300) FAC, PHI3
      GØ TØ 250
320  CONTINUE
      IF(IW.LE.1) GØ TØ 360
      IF(VVAR.EØ.VPAR) GØ TØ 340
      DØ 330 I=1, VPAR
330  EX(I)=0.E0
340  DØ 350 I=1, VVAR
350  EX(INX(I))=ANS(I)
      WRITE(IØUT,11600) (EX(I), I=1, VPAR)
360  CONTINUE
      GØ TØ 440
380  IF(IW.GE.0) WRITE(IØUT,12300)
      FAIL=2
      GØ TØ 450
390  IF(IW.GE.0) WRITE(IØUT,12000)
      FAIL=3
      RETURN
400  IF(IW.GE.0) WRITE(IØUT,11900) GAMA, LAM
      FAIL=4
      RETURN
410  IF(IW.GE.0) WRITE(IØUT,11400) I, (B(K), K=1, VPAR)
      FAIL=6
      RETURN
420  IF(IW.GE.0) WRITE(IØUT,10500)
      FAIL=5
      RETURN
430  IF(IW.GT.0) WRITE(IØUT,12400)
      FAIL=1
      GØ TØ 450
440  IF(ITER.LE.4) GØ TØ 80
C*****

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```

C     THE PROCEDURE HAS CONVERGED
C*****
      FAIL=0
      IF(IW.GE.0) WRITE(IOUT,11800)
      IF(IW.GE.0) WRITE(IOUT,11800)
      IF(IW.GE.0) WRITE(IOUT,12200) ITER
450   IF(IW.GE.0) WRITE(IOUT,12100) PHI,(B(I),I=1,NPAR)
C*****
C     CALCULATE THE STANDARD ERROR. USE NPTS-NVAR-NEXT AS THE DF.
C*****
      N0RM=SQRT(PHI/FL0AT(NPTS-NVAR-NEXT))
C*****
C     REMAKE THE PARTIAL MATRIX INSTEAD OF USING THE RE-SCALED MATRIX.
C     THIS COULD EASILY BE CHANGED.
C*****
      D0 460 I=1,NPAR
460   BDB(I)=0.E0
      D0 470 I=1,NVAR
      D0 470 J=I,NVAR
470   ATA(I,J)=0.E0
      D0 480 I=1,NPTS
      NPVFC=0
      CALL EYCALC(X,B,IE,I,EX,FV,IFXE,NPVFC,NC0MP,YN0RM)
      D0 480 L=1,NVAR
      J=INX(L)
      D0 480 M=L,NVAR
      K=INX(M)
480   ATA(L,M)=ATA(L,M)+EX(J)*EX(K)*WTS(I)
      CALL INVRT(NVAR,ATA,ID,IFU)
      IF(IFU.EQ.1) G0 T0 420
      D0 490 I=1,NVAR
      EX(I)=SQRT(ATA(I+1,I))
490   BDB(INX(I))=N0RM*EX(I)
C*****
C     CALCULATE THE CORRELATION MATRIX
C*****
      D0 500 I=1,NVAR
      D0 500 J=I,NVAR
500   ATA(I,J)=ATA(J+1,I)/(EX(I)*EX(J))
      IF(IW.LT.0) RETURN
C*****
C     PRINT THE PARAMETERS AND THE STANDARD ERRORS ASSOCIATED TO THEM
C*****
      WRITE(IOUT,11300) (INX(I),B(INX(I)),BDB(INX(I)),I=1,NVAR)
      IF(NVAR.EQ.1) G0 T0 520
      IF(FAIL.EQ.2) G0 T0 540
C*****
C     PRINT THE CORRELATION MATRIX
C*****
      WRITE(IOUT,11200)
      D0 510 I=1,NVAR

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510  WRITE(I0UT,11000) I,(ATA(J,I),J=1,I)
C*****
C    PRINT THE INVERSE MATRIX
C*****
      WRITE(I0UT,11800)
      WRITE(I0UT,11800)
520  WRITE(I0UT,11100)
      M=VVAR+1
      D0 530 I=2,M
      K=I-1
530  WRITE(I0UT,10400) I,(ATA(I,J),J=1,K)
540  CONTINUE
      WRITE(I0UT,11800)
      WRITE(I0UT,11800)
      WRITE(I0UT,10100)
      D0 550 I=1,VPTS
      VPVFC=1
      CALL EMCALC(X,B,IE,I,EX,FV,IFXE,VPVFC,VC0MP,YN0RM)
      SS=Y(I)-FV
550  WRITE(I0UT,10000) X(I),Y(I),FV,SS
      WRITE(I0UT,12500) YN0RM
      RETURN
10000 F0RMAT(4E14.8)
10100 F0RMAT(140,'FINAL DEVIATIONS'/1H,1X,'FREQUENCY',3X
      &,'0BS',13X,'CAL',12X,'0-C')
10200 F0RMAT(140,'INITIAL DEVIATIONS'/1H,13X,'0BS',13X,'CALC',12X,
      *'0-C')
10300 F0RMAT(140,'FAC=',E14.8,6X,' PHI(FAC) = ',E14.8)
10400 F0RMAT(14,'R0W ',I2/(1H,8E15.6))
10500 F0RMAT(140,'FINAL A TRANSP0SE A IS N0T P0SITIVE DEFINE')
10600 F0RMAT(140,' IN-L = ',E14.8,' PHI(IN-L) = ',E14.8,
      *' GAMA = ',F10.4)
10700 F0RMAT(140,' L/10 = ',E14.8,' PHI(L/10) = ',E14.8,
      *' GAMA = ',F10.4)
10800 F0RMAT(340L,7E14.8/(3H,16X,6E14.8))
10900 F0RMAT(140,' L*10 = ',E14.8,' PHI(L*10) V ',E14.8,
      *' GAMA = ',F10.4)
11000 F0RMAT(14,'R0W',I2/(1H,8E15.6))
11100 F0RMAT(140,' INVERSE MATRIX - LOWER TRIANGULAR P0RTI0N')
11200 F0RMAT(140,'C0RRELATI0N MATRIX LOWER TRIANGULAR P0RTI0N R0W BY R0W
      *PRINT')
11300 F0RMAT(140,' VARIABLE',6X,'PARAMETER VALUE',5X,'STANDARD ERR0R'/
      * (140,4X,I2,9X,E15.6,7X,E10.4))
11400 F0RMAT(140,'THE DIAG0NAL ELEMENT RESULTING FR0M THE PARTIAL WRT
      *B(',I2,') IS = 0.D'/1H,'THE P0INT AT WHICH THE FAILURE 0CCURED
      *IS'/(1H,7E14.8))
11500 F0RMAT(140,'ITERATI0N ',I3/1H,7E14.8/(1H,16X,6E14.8))
11600 F0RMAT(140,'DIR-VEC',5X,6E14.8/(140,12X,6E14.8))
11700 F0RMAT(141,T3,'NPAR',T11,'NVAR',T19,'NPTS',T27,'MAXITR',T37,'IW',
      *T42,'LINMIN',T56,'EP',T71,'DAC'/1H,T4,I2,T12,I2,T20,I2,T29,
      *I2,T37,I2,T44,I2,T53,E10.4,T63,E10.4)

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```
11800 FORMAT(1H0/)
11900 FORMAT(1H0,'GAMA = ',E14.8,' WHEN LAM = ',E14.8/
*1H , 'THERE PROBABLY EXISTS EXCESSIVELY HIGH CORRELATIONS BETWEEN
* THE PARAMETERS')
12000 FORMAT(1H0,'THE (ATA +LAM*I) MATRIX FAILED TO BE POS.DEF. ')
12100 FORMAT(1H0,7E14.8/(1H ,16X,6E14.8))
12200 FORMAT(1H0,'OPTIMAL POINT REACHED IN ',I5, ' ITERATIONS')
12300 FORMAT(1H0,'MAXIMUM NUMBER OF ITERATIONS REACHED-BEST POINT PRINTED')
12400 FORMAT(1H0,'THE DELTA-B VECTOR REDUCED TO CONVERGENCE LEVEL WHILE
*GAMA LESS THAN GAMA0. '/1H , 'THE POINT IS PROBABLY OPTIMAL WITHIN
*ROUNDING ERRORS. ')
12500 FORMAT(1X,'NORMALIZATION FACTOR = ',E14.8)
END
```

```

SUBROUTINE NEWLAM(ATA,LAM,BI,ATF,ANS,EX,GAMA,FVALUE,PHI,
&X,Y,WTS,B,IE,NVAR,NPTS,VPAR,IFAL,INX,IFXE,IR,NCOMP,YNORM)
C***** SUBROUTINE NEWLAM *****
C SUBROUTINE NEWLAM CALCULATES THE NEW LAMBDA FOR THE LEAST
C SQUARES PROCEDURE CONVERGENCE ACCORDING TO THE ALGORITHM
C STATED BY MARQUARDT (1963).
C
C
C ORIGINAL PROGRAMER JORGE PARRA
C MODIFIED BY JEFF DANIELS
C*****
DIMENSION ATA(7,7),BI(7),ATF(7),ANS(7),EX(7)
DIMENSION Y(IE),X(IE),WTS(IE),B(7)
DIMENSION INX(20),IFXE(7)
REAL LAM
C0V=57.295779E0
IFAL=0
D0 10 I=1,NVAR
ATA(I,I)=1.E0
10 ATA(I,I)=ATA(I,I)+LAM
ERRORX=.001E+0
CALL DS3L (ATA,ATF,ANS,25,NVAR,ERRORX,10,IDET,1,IFAIL)
1005 FORMAT(I5)
1006 FORMAT(3(4X,E14.8))
IF(IFAIL.EQ.1) G0 T0 20
IFAL=1
RETURN
20 D0 30 I=1,NVAR
30 ANS(I)=ANS(I)/EX(I)
IFAIL=0
IF(NVAR.NE.1) G0 T0 40
GAMA=0.E0
G0 T0 70
40 SUM1=0.E0
SUM2=0.E0
SUM3=0.E0
D0 50 I=1,NVAR
SUM1=SUM1+ANS(I)*ATF(I)
SUM2=SUM2+ATF(I)**2
50 SUM3=SUM3+ANS(I)**2
C0SGAM=SUM1/SQRT(SUM2*SUM3)
UV=SQRT(1.E0-C0SGAM*C0SGAM)
GAMA=ATAN(UV/C0SGAM)*C0V
IF(C0SGAM.GT.0.E0) G0 T0 60
GAMA=180.E0-GAMA
IF(LAM.LT.1.E0) G0 T0 60
IFAIL=1
60 IF(IFAIL.EQ.0) G0 T0 70
IFAL=2
RETURN
70 D0 80 I=1,NVAR
K=INX(I)
80 BI(K)=B(K)+ANS(I)
PHI=0.E0
D0 90 I=1,NPTS
NPVFC=1
CALL EMCALC(X,BI,IE,I,EX,FV,IFXE,NPVFC,NCOMP,YNORM)
90 PHI=PHI+(Y(I)-FV)**2*WTS(I)
RETURN
END

```

SUBROUTINE INVRT(N,A,IE,IFL)

C\*\*\*\*\*SUBROUTINE INVRT\*\*\*\*\*  
 C THIS PROGRAM PERFORMS A MATRIX INVERSION ON THE  
 C TWO DIMENSIONAL ARRAY A.  
 C\*\*\*\*\*

```

      DIMENSION A(7,7)
      IFL=0
      DO 40 I=1,N
      I1=I+1
      DO 40 J=I,N
      J1=J+1
      X=A(I,J)
      K=I-1
10     IF(K.LT.1) GO TO 20
      X=X-A(J1,K)*A(I1,K)
      K=K-1
      GO TO 10
20     IF(J.NE.I) GO TO 30
      IF(X.LE.0.E0) GO TO 90
      Y=1.E0/SQRT(X)
      A(I1,I)=Y
      GO TO 40
30     A(J1,I)=X*Y
40     CONTINUE
      NL=N-1
      DO 60 I=1,NL
      KL=I+1
      DO 60 J=KL,N
      Z=0.E0
      J1=J+1
      K=J-1
50     IF(K.LT.I) GO TO 60
      Z=Z-A(J1,K)*A(K+1,I)
      K=K-1
      GO TO 50
60     A(J1,I)=Z*A(J1,J)
      J1=N+1
      DO 80 I=1,N
      DO 80 J=I,N
      Z=0.D0
      KL=J+1
      DO 70 K=KL,J1
70     Z=Z+A(K,J)*A(K,I)
80     A(KL,I)=Z
      RETURN
90     IFL=1
      RETURN
      END

```

SUBROUTINE DSOL (A,B,X,NAMAX,NA,ERRORX,MAXIT,K,IENTRY,ISING)

THIS SUBROUTINE SOLVES A SET OF SIMULTANEOUS LINEAR EQUATIONS USING CROUTS FACTORIZATION WITH PARTIAL PIVOTING, EQUILIBRATION, ITERATIVE IMPROVEMENT, AND DOUBLE PRECISION OF PRODUCTS TO REDUCE ROUND-OFF ERROR.

SUBROUTINE PARAMETERS .....

A = THE COEFFICIENT MATRIX  
 B = THE CONSTANT VECTOR  
 X = THE SOLUTION VECTOR (INFORMATION RETURNED)  
 NAMAX = THE NUMBER OF ROWS IN THE DIMENSION OF A IN THE CALLING PROGRAM.  
 NA = THE NUMBER OF ROWS IN A  
 ERRORX = THE MAXIMUM ALLOWABLE ERROR BETWEEN SUCCESSIVE ITERATIONS  
 MAXIT = THE MAXIMUM ALLOWABLE NUMBER OF ITERATIONS  
 K = THE ACTUAL NUMBER OF ITERATIONS PERFORMED (INFORMATION RETURNED)  
 IENTRY = 1 (FIRST CALL TO SUBROUTINE)  
           = 2 (SUBSEQUENT CALLS WITH UNCHANGED A)  
 ISING = 1 THE MATRIX IS NON-SINGULAR (INFORMATION RETURNED)  
           = 2 THE MATRIX IS SINGULAR. PROCEDURE DISCONTINUED.

.....  
 IF A SINGULAR MATRIX IS ENCOUNTERED THE PROCEDURE IS DISCONTINUED AND ERROR MESSAGES ARE PRINTED OUT.  
 IF A SOLUTION IS REQUIRED FOR THE SAME A-MATRIX, BUT FOR DIFFERENT B-VECTORS, THE SUBROUTINE MAY BE SUBSEQUENTLY DIRECTED TO ENTRY POINT -- TWO.

.....  
 PROGRAMED BY  
 DONALD SNYDER 1968  
 COLORADO SCHOOL OF MINES

DOUBLE PRECISION S  
 DIMENSION A(7,7),B(7),X(7),BIGB(7,7),Q(7),R(7)  
 DIMENSION IPIV(10)  
 IOUT=2

GO TO (100,200),IENTRY

ENTRY -- ONE .....

100 CONTINUE

STORE MATRIX A IN ARRAY BIGB

DO 101 I=1,NA

DO 101 J=1,NA

101 BIGB(I,J)=A(I,J)

DETERMINE ELEMENT WITH MAXIMUM MODULUS IN EACH ROW. STORE IN Q.

DO 103 I=1,NA

PLAINQ=0.0

```

      D3 102 J=1,NA
      TEST=ABS(A(I,J))
      IF (TEST.GT.PLAINQ) PLAINQ=TEST
102  CONTINUE
      Q(I)=PLAINQ
C     BOX 13 -- CHECK FOR SINGULARITY .....
      IF (PLAINQ.LT.0.1E-30) G0 T0 301
C
C     DETERMINE THE ELEMENTS OF THE LOWER TRIANGULAR FACTOR OF A
C     STORE OVER BIGB
      D3 110 IR=1,NA
      AMAX=0.0
      IRMAX=IR
      LIM=IR-1
      D3 106 I=IR,NA
      S=DBLE(BIGB(I,IR))
      IF (IR.EQ.1) G0 T0 105
      D3 104 J=1,LIM
104  S=S-BIGB(I,J)*BIGB(J,IR)
105  BIGB(I,IR)=SNGL(S)
      TEST=ABS(BIGB(I,IR)/Q(I))
      IF (TEST.LE.AMAX) G0 T0 106
      AMAX=TEST
      IRMAX=I
106  CONTINUE
C
C     BOX 37 - CHECK FOR SINGULARITY .....
      IF (AMAX.LT.0.1E-30) G0 T0 302
C
      Q(IRMAX)=Q(IR)
      IPIV(IR)=IRMAX
C
C     INTERCHANGE THE IR AND IRMAX ROWS OF BIGB
      D3 107 I=1,NA
      PLAINQ=BIGB(IR,I)
      BIGB(IR,I)=BIGB(IRMAX,I)
107  BIGB(IRMAX,I)=PLAINQ
C
C     DETERMINE THE ELEMENTS OF THE UPPER TRIANGULAR FACTOR OF A
C     STORE OVER BIGB
      IF (IR.EQ.NA) G0 T0 110
      LOW=IR+1
      D3 109 I=LOW,NA
      S=DBLE(BIGB(IR,I))
      IF (IR.EQ.1) G0 T0 109
      D3 108 J=1,LIM
108  S=S-BIGB(IR,J)*BIGB(J,I)
109  BIGB(IR,I)=SNGL(S/DBLE(BIGB(IR,IR)))

```

```

110 CONTINUE
C
C   FACTORIZATION IS COMPLETE. NOW FIND SOLUTION
C
C   ENTRY -- TWO .....
200 CONTINUE
C
C   PRESET ITERATION PARAMETERS, AND INITIALIZE SOLUTION AND ERROR VEC
K=0
IFIN=0
ISING=1
DØ 201 I=1,NA
X(I)=0.0
201 R(I)=B(I)
- 202 CONTINUE
C
C   PERFORM FORWARDS SUBSTITUTION
DØ 204 I=1,NA
L=IPIV(I)
S=DBLE(R(L))
R(L)=R(I)
IF (I.EQ.1) GØ TØ 204
LIM =I-1
DØ 203 J=1,LIM
203 S=S-BIGB(I,J)*R(J)
204 R(I)=SGL(S/DBLE(BIGB(I,I)))
C
C   PERFORM BACKWARDS SUBSTITUTION
I=NA+1
DØ 206 N=1,NA
I=I-1
S=DBLE(R(I))
IF (I.EQ.NA) GØ TØ 206
LØW=I+1
DØ 205 J=LØW,NA
205 S=S-BIGB(I,J)*R(J)
206 R(I)=S
C
C   COMPUTE NORMS
AVØRMX=0.0
AVØRME=0.0
IF (K.EQ.0) GØ TØ 210
DØ 207 I=1,NA
TEST=ABS(X(I))
IF (TEST.GT.AVØRMX) AVØRMX=TEST
TEST=ABS(R(I))
IF (TEST.GT.AVØRME) AVØRME=TEST
207 CONTINUE
C
C   IF (K.NE.1) GØ TØ 208
C

```

```

C      BØX 94 - CHECK FØR SINGULARITY .....
      IF ((ANØRME/ANØRMX).GT.0.5) GØ TØ 303
C
209 IF ((ANØRME/ANØRMX).LT.ERRØRX) IFIN=1
210 CØNTINUE
C
C      CØMPUTE CURRENT APPROXIMATE SØLUTIØN
      DØ 211 I=1,NA
211 X(I)=X(I)+R(I)
C
C      CØMPUTE ERRØR AND STØRE BACK IN R
      DØ 213 I=1,NA
      S=DBLE(B(I))
      DØ 212 J=1,NA
212 S=S-A(I,J)*X(J)
213 R(I)=SNGL(S)
      K=K+1
      IF (IFIN.EØ.1) RETURN
      IF (K.LT.MAXIT) GØ TØ 202
C
C      PRINT ERRØR MESSAGE AND RETURN
      WRITE(IØUT,11) MAXITR
      RETURN
301 IBØX=13
      GØ TØ 304
302 IBØX=37
      GØ TØ 304
303 IBØX=94
304 WRITE(IØUT,12) IBØX
      ISING=2
C
11 FØRMAT(140,53HSØLUTIØN HAS NOT CØNVERGED TØ DESIRED ACCURACY AFTER
1,15,11H ITERATIØNS)
12 FØRMAT(140,31HSINGULAR MATRIX DETECTED AT BØX,13)
C
      RETURN
      END

```

```

SUBROUTINE EMCALC(FF, B, IE, I, DERIV, FV, IFXE, VPVFC, NCOMP, YNORM)
C***** SUBROUTINE EMCALC *****
C THIS SUBROUTINE HANDLES BOTH THE CALCULATION OF THE FORWARD
C SOLUTION AND THE PARTIAL DERIVATIVES OF THE FORWARD
C SOLUTION WITH RESPECT TO THE LAYERED EARTH PARAMETERS
C AND THE SOURCE RECEIVER SEPERATION.
C
C***** VARIABLES*****
C FF= FREQUENCY IN HERTZ
C F= ANGULAR FREQUENCY
C D(I)= NORMALIZED THICKNESS FOR THE I-TH LAYER
C RK(I)= NORMALIZED RESISTIVITIES FOR THE I-TH LAYER
C DEL= MODIFIED WAVE NUMBER
C
C*****
      COMPLEX H, E
      COMMON /PD/ EXX(6), HX(6), DIVX(6), KFXE(6), EC, HZ, ECR, ECI, HZR, HZI
&, EXR(6), EXI(6), HXR(6), HXI(6)
      COMMON /QC/DD(3), RKK(3), RHH(3), V, DEL, RR
      COMMON /CV/CI, DS, ANG, F, XX, TM
      COMMON /PART/ RK(3), RH(3), HP(3), HPP(3), D(3),
&X, R, VB
      DIMENSION B(7), FF(25), DERIV(7), IFXE(6)

      R=B(1)
      VB=2*V
      D0 1 JCT=1, V
      VCT=2*JCT
      RH(JCT)=B(VCT)
      HP(JCT)=B(VCT+1)
1     CONTINUE
      D0 3 JC=1, VB
3     KFXE(JC)=IFXE(JC)
      TM=12.566371E-7
      F=FF(1)*6.2831853
      DEL=SQRT(2.*RH(1)/(TM*F))
      D0 2 JJ=1, V
2     D(JJ)=2.*HP(JJ)/DEL
      RK(JJ)=RH(1)/RH(JJ)
      X=ALOG(R/DEL)
      P=1.001
      RR=R
      XX=X
      D0 9 J=1, V
9     RHH(J)=RH(J)
      RKK(J)=RK(J)
      DD(J)=D(J)
      CALL FVAL(E, H, NCOMP)
      EC=CABS(E)
      HZ=CABS(H)
      ECR=REAL(E)

```

```

ECI=AIMAG(E)
HZR=REAL(H)
HZI=AIMAG(H)
IF(NPVFC.EQ.0) GO TO 17
IF(NCOMP.EQ.10) FV=EC/YNORM
IF(NCOMP.EQ.40) FV=EC/YNORM
IF(NCOMP.EQ.70) FV=EC/YNORM
IF(NCOMP.EQ.20) FV=ECR/YNORM
IF(NCOMP.EQ.50) FV=ECR/YNORM
IF(NCOMP.EQ.80) FV=ECR/YNORM
IF(NCOMP.EQ.30) FV=ECI/YNORM
IF(NCOMP.EQ.60) FV=ECI/YNORM
IF(NCOMP.EQ.90) FV=ECI/YNORM
IF(NCOMP.EQ.11) FV=HZ/YNORM
IF(NCOMP.EQ.21) FV=HZR/YNORM
IF(NCOMP.EQ.31) FV=HZI/YNORM
17 CONTINUE
IF(NPVFC.EQ.1) GO TO 4
CALL PARDER (IFXE, DERIV, NCOMP, YNORM)
4 CONTINUE
RETURN
END

```

```

SUBROUTINE PARDER(IFXE, DERIV, VCOMP, VNORM)
C*****SUBROUTINE PARDER*****
C THIS SUBROUTINE CONTROLS THE PARTIAL DIFFERENTIATION WRT THE
C LAYERING PARAMETERS (RESISTIVITY AND THICKNESS)
C*****
COMMON /PD/ EXX(6),HX(6),DIVX(6),IFXE(6),EC,HZ,ECR,ECI,
&HZR,HZI,EXR(6),EXI(6),HXR(6),HXI(6)
COMMON /DC/ DD(3),RKK(3),RHH(3),N,DEL,RR
COMMON /CV/ CI,DS,AVG,F,XX,TM
COMMON /PART/ RK(3),RH(3),HP(3),HPP(3),
&DC(3),X,R,NB
DIMENSION DERIV(7),IFXE(6)
COMPLEX H,E
JP=1
IF(IFXE(JP).EQ.0) GO TO 6
P=1.001
RR=P*RR
XX=ALOG(RR/DEL)
DIVX(JP)=RR-R
CALL FVAL(E,H,VCOMP)
EXX(JP)=CABS(E)
HX(JP)=CABS(H)
EXR(JP)=REAL(E)
EXI(JP)=AIMAG(E)
HXR(JP)=REAL(H)
HXI(JP)=AIMAG(H)
XX=X
GO TO 7
6 EXX(JP)=EC
HX(JP)=HZ
EXR(JP)=ECR
EXI(JP)=ECI
HXR(JP)=HZR
HXI(JP)=HZI
7 JP=JP+1
DO 12 J=1,N
IF(IFXE(JP).EQ.0) GO TO 8
RHH(J)=P*RH(J)
RKK(J)=RHH(1)/RHH(J)
DIVX(JP)=RHH(J)-RH(J)
CALL FVAL(E,H,VCOMP)
EXX(JP)=CABS(E)
HX(JP)=CABS(H)
EXR(JP)=REAL(E)
EXI(JP)=AIMAG(E)
HXR(JP)=REAL(H)
HXI(JP)=AIMAG(H)
RHH(J)=RH(J)
RKK(J)=RK(J)
GO TO 13
8 EXX(JP)=EC

```

```

HX(JP)=HZ
EXR(JP)=ECR
EXI(JP)=ECI
HXR(JP)=HZR
HXI(JP)=HZI
13  JP=JP+1
    IF(J.EQ.V) GO TO 12
    IF(IFXE(JP).EQ.0) GO TO 10
    HPP(J)=P*HP(J)
    DD(J)=2.*HPP(J)/DEL
    DIX(JP)=HPP(J)-HP(J)
    CALL FVAL(E,H,VCOMP)
    EXX(JP)=CABS(E)
    HX(JP)=CABS(H)
    EXR(JP)=REAL(E)
    EXI(JP)=AIMAG(E)
    HXR(JP)=REAL(H)
    HXI(JP)=AIMAG(H)
    HPP(J)=HP(J)
    DD(J)=2.*HP(J)/DEL
    GO TO 12
10  EXX(JP)=EC
    HX(JP)=HZ
    EXR(JP)=ECR
    EXI(JP)=ECI
    HXR(JP)=HZR
    HXI(JP)=HZI
12  JP=JP+1
5   CONTINUE
    CALL FINDIF(NB,VCOMP,DERIV)
    DO 19 JPZ=1,NB
19  DERIV(JPZ)=DERIV(JPZ)/YNORM
4   CONTINUE
    RETURN
    END

```

```

SUBROUTINE FINDIF (NB,NCOMP,DERIV)
C*****SUBROUTINE FINDIF*****
C THIS SUBROUTINE CALCULATES THE PARTIAL DERIVATIVES OF THE
C PARAMETERS USING FINITE DIFFERENCES
C*****

      COMMON /PD/ EX(6),HX(6),DINX(6),INX(6),EC,HZ,ECR,ECI,HZR,HZI,
&EXR(6),EXI(6),HXR(6),HXI(6)
      DIMENSION DERIV(7)
      DO 1 JP=1,NB
      IF(INX(JP).EQ.0) GO TO 2
      IF(NCOMP.EQ.10) GO TO 10
      IF(NCOMP.EQ.40) GO TO 10
      IF(NCOMP.EQ.70) GO TO 10
      GO TO 3
10     DERIV(JP)=(EX(JP)-EC)/DINX(JP)
1005   FORMAT(3(2X,E14.8))
      GO TO 1
      3   IF(NCOMP.NE.11) GO TO 4
      DERIV(JP)=(HX(JP)-HZ)/DINX(JP)
      GO TO 1
      4   IF(NCOMP.EQ.20) GO TO 11
      IF(NCOMP.EQ.50) GO TO 11
      IF(NCOMP.EQ.80) GO TO 11
      GO TO 5
11     DERIV(JP)=(EXR(JP)-ECR)/DINX(JP)
      GO TO 1
      5   IF(NCOMP.EQ.30) GO TO 12
      IF(NCOMP.EQ.60) GO TO 12
      IF(NCOMP.EQ.90) GO TO 12
      GO TO 6
12     DERIV(JP)=(EXI(JP)-ECI)/DINX(JP)
      GO TO 1
      6   IF(NCOMP.NE.21) GO TO 7
      DERIV(JP)=(HXR(JP)-HZR)/DINX(JP)
      GO TO 1
      7   IF(NCOMP.NE.31) GO TO 1
      DERIV(JP)=(HXI(JP)-HZI)/DINX(JP)
      GO TO 1
      2   DERIV(JP)=0.0
      1   CONTINUE
      RETURN
      END

```

## SUBROUTINE FVAL (EC,H,VCOMP)

```

*****SUBROUTINE FVAL *****
C THIS SUBROUTINE CALCULATES THE FORWARD SOLUTION ON COMMAND FROM
C EMCALC USING LINEAR DIGITAL FILTER THEORY
C
*****VARIABLES*****
C CH= COEFFICIENTS FOR CALCULATING THE J1 HANKEL TRANSFORM
C YH= ABCISSA VALUES FOR CH
C CE= COEFFICIENT FOR CALCULATING THE J0 HANKEL TRANSFORM
C YE= ABCISSA VALUES FOR CALCULATING CE
C SUM= CONVOLUTION SUM FOR HZ
C SUEY1, SUEX1, SUEY0, SUEX0, ARE AS DEFINED ON PAGE 10 OF THESIS
C SUEY1= CONVOLUTION SUM FOR EY1
C SUEX1= CONVOLUTION SUM FOR EX3
C SUEY0= CONVOLUTION SUM FOR EX1
C SUEX0= CONVOLUTION SUM FOR EX2
C H1= CLOSED FORM HOMOGENEOUS HALF-SPACE EXPRESSION FOR HZ
C EX1= CLOSED FORM EXPRESSION FOR A HOMOGENEOUS HALF-SPACE
C      ( X-COMPONENT)
C EY1= CLOSED FORM EXPRESSION FOR A HOMOGENEOUS HALF-SPACE
C      (Y-COMPONENT)
C EX= X-COMPONENT OF THE ELECTRIC FIELD FOR A LAYERED EARTH
C EY= Y-COMPONENT OF THE ELECTRIC FIELD FOR A LAYERED EARTH
C EC= COMBINED ELECTRIC FIELD FOR A LAYERED EARTH
C F1= LAYERED EARTH CORRECTION FACTOR
C Z1= LAYERED EARTH CORRECTION FACTOR
C V1= NORMALIZED FIRST LAYER PSEUDO WAVE NUMBER
C
C
C*****

```

```

COMMON /CV/ CI, DS, ANG, F, X, TM
COMMON /QC/ D(3), RK(3), RH(3), N, DEL, R
COMPLEX EX1, EX0, EY1, EY0, SUEX1, SUEX0, SUEY1, SUEY0, EX, EY
COMPLEX H, V1, F1, SUM, Z1, E1, EC, H1, C1, C2
DIMENSION YH(48), CH(48), YE(61), CE(61)

```

DATA

```

&YH/-4.5307316E 0, -4.3004731E 0, -4.0702146E 0, -3.8399561E 0,
& -3.6096976E 0, -3.3794391E 0, -3.1491806E 0, -2.9189221E 0,
& -2.6886636E 0, -2.4584051E 0, -2.2281466E 0, -1.9978881E 0,
& -1.7676296E 0, -1.5373711E 0, -1.3071126E 0, -1.0768541E 0,
& -8.4659563E-1, -6.1633713E-1, -3.8607863E-1, -1.5582013E-1,
& 7.4438369E-2, 3.0469687E-1, 5.3495537E-1, 7.6521387E-1,
& 9.9547237E-1, 1.2257309E 0, 1.4559394E 0, 1.6862479E 0,
& 1.9165064E 0, 2.1467649E 0, 2.3770234E 0, 2.6072819E 0,
& 2.8375404E 0, 3.0677939E 0, 3.2980574E 0, 3.5233159E 0,
& 3.7585744E 0, 3.9888329E 0, 4.2190914E 0, 4.4493499E 0,
& 4.6796084E 0, 4.9098669E 0, 5.1401254E 0, 5.3703839E 0,
& 5.6006424E 0, 5.8309009E 0, 6.0611594E 0, 6.2914179E 0/

```

DATA

```

&CH/3.1010561E-6, 1.8802098E-5, 5.4819540E-5, 9.2891602E-6,
& 1.5523239E-4, 3.0344652E-5, 3.5338744E-4, 1.4793002E-4,

```

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& 7.7342377E-4, 5.3570857E-4, 1.7170605E-3, 1.6387239E-3,
& 3.9247683E-3, 4.5796508E-3, 9.2111468E-3, 1.2130467E-2,
& 2.1938415E-2, 3.0853660E-2, 5.1973594E-2, 7.4661566E-2,
& 1.1775455E-1, 1.6353574E-1, 2.3127545E-1, 2.7368461E-1,
& 2.8059285E-1, 1.2875840E-1, -1.5380437E-1, -4.5659951E-1,
& -3.6077766E-2, 4.2985683E-1, -2.1506075E-1, -2.3624312E-2,
& 8.9316746E-2, -7.4344203E-2, 4.8572965E-2, -3.0083372E-2,
& 1.8846544E-2, -1.2158687E-2, 8.0708759E-3, -5.4706275E-3,
& 3.7554604E-3, -2.5929707E-3, 1.7909426E-3, -1.2320277E-3,
& 8.4095286E-4, -5.6749747E-4, 3.7718405E-4, -1.5891335E-4/

```

## DATA

```

&YE/-6.8348046E 0, -6.6045461E 0, -6.3742876E 0, -6.1440291E 0,
& -5.9137706E 0, -5.6835121E 0, -5.4532536E 0, -5.2229951E 0,
& -4.9927366E 0, -4.7624781E 0, -4.5322196E 0, -4.3019611E 0,
& -4.0717026E 0, -3.8414441E 0, -3.6111856E 0, -3.3809271E 0,
& -3.1506686E 0, -2.9204101E 0, -2.6901516E 0, -2.4598931E 0,
& -2.2296346E 0, -1.9993761E 0, -1.7691176E 0, -1.5388591E 0,
& -1.3086006E 0, -1.0783421E 0, -8.4803358E-1, -6.1782508E-1,
& -3.8756658E-1, -1.5730808E-1, 7.2950416E-2, 3.0320892E-1,
& 5.3346742E-1, 7.6372592E-1, 9.9398442E-1, 1.2242429E 0,
& 1.4545014E 0, 1.6847599E 0, 1.9150184E 0, 2.1452769E 0,
& 2.3755354E 0, 2.6057939E 0, 2.8360524E 0, 3.0663109E 0,
& 3.2965694E 0, 3.5268279E 0, 3.7570864E 0, 3.9873449E 0,
& 4.2176034E 0, 4.4478619E 0, 4.6731204E 0, 4.9083789E 0,
& 5.1386374E 0, 5.3688959E 0, 5.5991544E 0, 5.8294129E 0,
& 6.0596714E 0, 6.2899299E 0, 6.5201884E 0, 6.7504469E 0,
& 6.9807054E 0/

```

## DATA

```

&CE/7.3260937E-4, 5.6326423E-4, 1.3727237E-4, 7.5331222E-4,
& 3.5918326E-4, 1.0500608E-3, 7.1530982E-4, 1.5160070E-3,
& 1.2841617E-3, 2.2497985E-3, 2.1906186E-3, 3.4076782E-3,
& 3.6321245E-3, 5.2376028E-3, 5.9212519E-3, 8.1315877E-3,
& 9.5527062E-3, 1.2708615E-2, 1.5305589E-2, 1.9941086E-2,
& 2.4396626E-2, 3.1333652E-2, 3.8683065E-2, 4.9127993E-2,
& 6.0824806E-2, 7.6314344E-2, 9.3928346E-2, 1.1545027E-1,
& 1.3868663E-1, 1.6248847E-1, 1.8114332E-1, 1.8424433E-1,
& 1.5556741E-1, 6.8592481E-2, -8.8339029E-2, -2.8819226E-1,
& -3.5565260E-1, -5.6288677E-2, 4.8186942E-1, -5.1516453E-2,
& -2.6102989E-1, 2.1416490E-1, -9.4490687E-2, 2.6196370E-2,
& -5.1097828E-4, -6.6032948E-3, 7.5193619E-3, -6.7854344E-3,
& 5.8044372E-3, -4.9354894E-3, 4.2323106E-3, -3.6733648E-3,
& 3.2266260E-3, -2.8649137E-3, 2.5677680E-3, -2.3202655E-3,
& 2.1115187E-3, -1.9334662E-3, 1.7800248E-3, -1.6465436E-3,
& 1.3468317E-3/

```

SUEX1=CMPLX(0.0,0.0)

SUEY1=CMPLX(0.0,0.0)

SUEX0=CMPLX(0.0,0.0)

SUEY0=CMPLX(0.0,0.0)

SUM=CMPLX(0.0,0.0)

BB=EXP(X)

CI=CMPLX(0.0,1.0)

```

      C2=CMPLX(1.0,1.0)
C*****
CALCULATE THE FIRST ORDER HANKEL TRANSFORMS
C*****
      DØ 8 J=1,48
      Y=EXP(-(X-YH(J)))
      CALL CALC(F1,Z1,V1,Y)
      EX1=C1*V1*(Z1-1)+2.*V1*(1-F1)/((Y+V1*F1)*(Y+V1))
      SUEX1=SUEX1+CH(J)*EX1
      H=Y*Y*V1*(1.-F1)
      H=H/((Y+V1*F1)*(Y+V1))
      SUM=SUM+CH(J)*H
9      CONTINUE
      SUEY1=SUEX1
      H=(1/BB)*C1*DS*SIN(ANG)*SUM/(6.2831853*DEL*DEL)
      R2=R*R
      H1=3.- (3.+3.0*C2*BB+2.*C1*BB**2)*CEXP(-C2*BB)
      H1=(-C1*C1*DS*DEL*DEL*SIN(ANG)*H1)/(12.566371*R2*R2)
12     FORMAT(2(1X,E14.8))
      H=H+H1
      IF(MØD(NCØMP,10).EQ.1) GØ TO 6
      SUM=CMPLX(0.0,0.0)
C*****
CALCULATE THE ZERO TH ORDER HANKEL TRANSFORMS
C*****
      DØ 4 I=1,61
      Y=EXP(-(X-YE(I)))
      CALL CALC(F1,Z1,V1,Y)
      EY0=Y*(C1*V1*(Z1-1.)+2.*V1*(1.-F1)/((Y+V1)*(Y+V1*F1)))
      EX0=(CØS(ANG)**2/DEL)*EY0-(1/DEL)*2.*Y*V1*(1.-F1)/
      &((Y+V1*F1)*(Y+V1))
      SUEX0=SUEX0+CE(I)*EX0
      SUEY0=SUEY0+CE(I)*EY0
      EC=C1*V1*Y*(Z1-1.0)/2.0
      EC=EC-Y*V1*(1.0-F1)/((Y+V1*F1)*(Y+V1))
4      SUM=SUM+CE(I)*EC
      EC=C1*TM*F*C1*DS*SIN(2.*ANG)*SUM/(12.566371*DEL*BB)
      E1=-.500+(1+C2*BB)*CEXP(-C2*BB)
      E1=C1*DS*RH(1)*SIN(2.*ANG)*E1/(6.2831853*R*R2)
      EC=E1+EC
      EX=(C1*TM*F*C1*DS/(12.566371*BB))*(SUEX0-CØS(2.*ANG)
      &*SUEX1/R)
      EY=(C1*TM*F*C1*DS/(12.566371*BB))*SIN(2.*ANG)*(-SUEY1/R
      &+SUEY0/(2.*DEL))
      EX1=C1*DS*RH(1)*((3*CØS(ANG)**2-2)+(1.+(1+C1)*BB)*CEXP(-C2*BB)
      &)/(6.2831853*R*R2)
      EY1=3.*RH(1)*C1*DS*SIN(ANG)*CØS(ANG)/(6.2831853*R*R2)
      EX=EX+EX1
      EY=EY+EY1
      IF(NCØMP.GE.10.AND.NCØMP.LE.30) GØ TO 7
      IF(NCØMP.GE.40.AND.NCØMP.LE.60) GØ TO 11

```

```
IF(NCOMP.GE.70.AND.NCOMP.LE.90) GO TO 9
7  EC=EC
   GO TO 6
11 EC=EX
   GO TO 6
9  EC=EY
6  CONTINUE
   RETURN
   END
```

```

SUBROUTINE CALC(F1,Z1,V1,Y)
C*****
C SUBROUTINE CALC CALCULATES THE LAYERED EARTH CORRECTION FACTORS
C (F1 AND Z1) FOR UP TO 3 LAYERS*****
COMMON /QC/ D(3),RK(3),RH(3),N,DEL,R
COMPLEX V1,F1,V2,AEX,Z1
Y2=Y*Y
T=2.*RK(N)
V2=CSQRT(CMPLX(Y2,T))
DO 9 LL=2,N
IF(LL.GT.2) GO TO 4
F1=CMPLX(1.0,0.0)
Z1=CMPLX(1.0,0.0)
4 I=N-LL+1
DD=D(I)
T=2.*RK(I)
V1=CSQRT(CMPLX(Y2,T))
AEX=(1.-CEXP(-V1*DD))/(1.+CEXP(-V1*DD))
F1=(V2*F1+V1*AEX)/(V1+V2*F1*AEX)
Z1=(V2*RH(I+1)*Z1+V1*RH(I)*AEX)/(V1*RH(I)+V2*RH(I+1)*Z1*AEX)
9 V2=V1
RETURN
END

```

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